

# Physical properties of semiconductor nanostructures: theory and simulations

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*PULSE school “Epitaxy bases and promises”  
Porquerolles, France  
July 2-8, 2021*

# STEM group @ LPS, Orsay



**Group Leader: O. Stéphan, LPS, Université Paris Saclay**  
Electron microscopy group, specialized in ultra-high resolution EELS

<https://www.stem.lps.u-psud.fr>

# Outline

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- › Introduction on nanoscience
- › Ab initio simulations: Density Functional Theory (DFT)
- › Quantum confinement in Si nanowires (Si NWs)
  - Electronic structure
  - Transport properties
  - Doping effects
- › Crystal phase engineering in Si NWs
  - Crystal phase effects on bulk
  - Crystal phase effects on NWs
  - Homojunctions

## › Conclusions

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- Electronic structure
- Transport properties
- Doping effects

## Crystal phase engineering in Si NWs

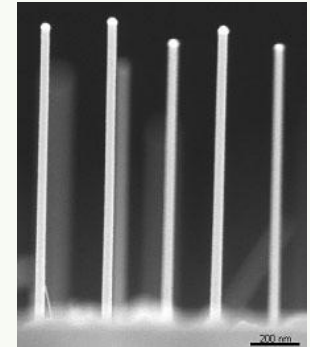
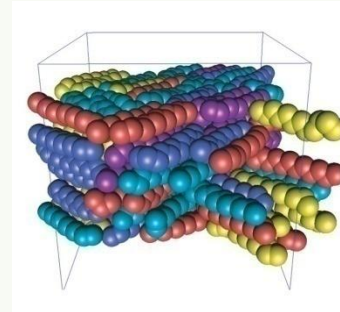
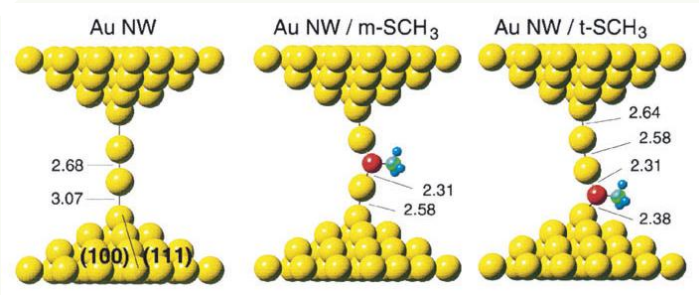
- Crystal phase effects on bulk
- Crystal phase effects on NWs
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## Conclusions

# Nanoscience

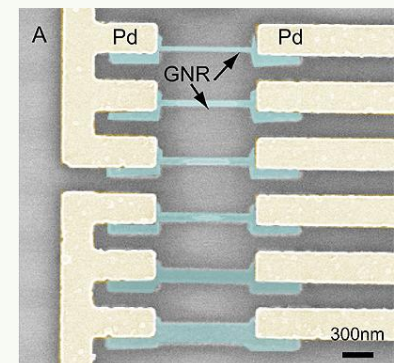
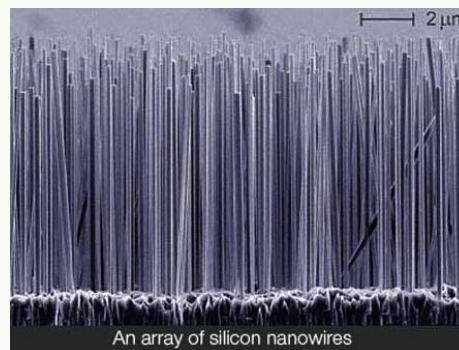
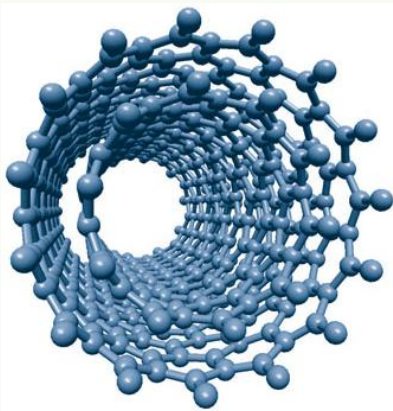


# Nanoscience



The study of structures, dynamics, and properties of systems in which one or more of the spatial dimensions is nanoscopic (1-100 nm) Dynamics and properties that are distinctly different (often in an extraordinary way) from both small-molecule systems and systems macroscopic in all dimensions

*U.S. National Nanotechnology Initiative, nano.gov*



# The importance of size



# The importance of size

Gold bar



Gold coins



Gold flakes



Gold nanoparticles



*As the size of the material is reduced, and the nanoscale regime is reached, it is possible that the same material will display totally different properties*

# Nano building blocks

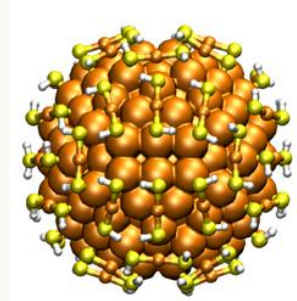
# Nano building blocks

*What are the nano building blocks that would play an analogous role to macro building blocks ?*

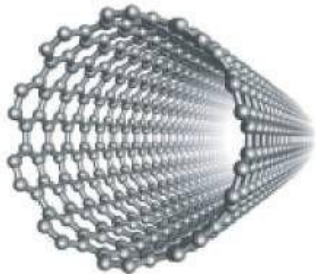
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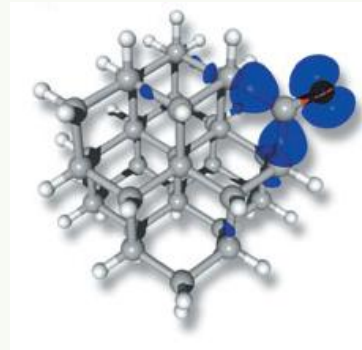
Clusters and molecular nanostructures



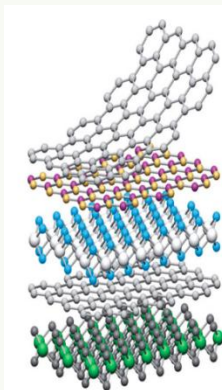
Nanotubes and related systems



Quantum wells, wires, films and dots



2D materials



# From nanoscience to ...

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Chemistry

Physics

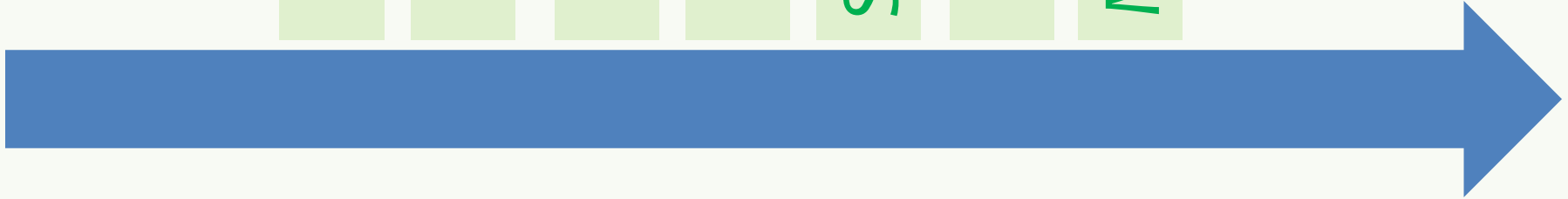
Biology

Engineering

Surface Science

Biotechnology

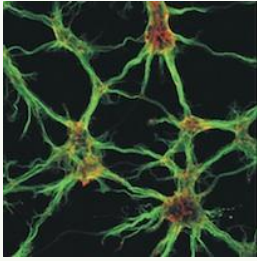
Molecular Biology



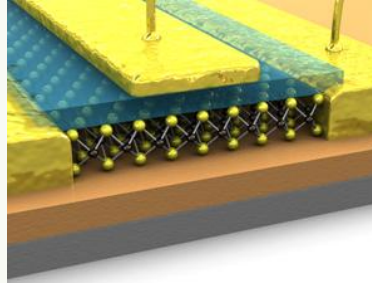
# ... nanotechnology



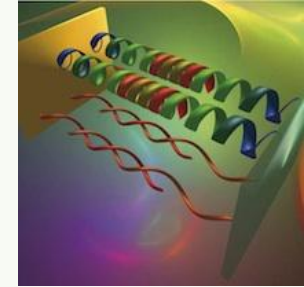
# ... nanotechnology



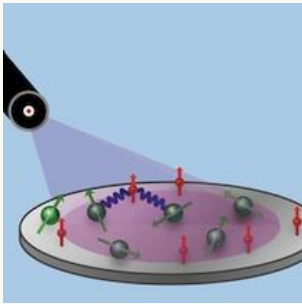
Biomedicine



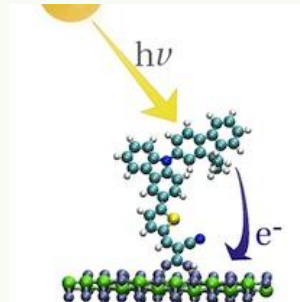
Nanoelectronics



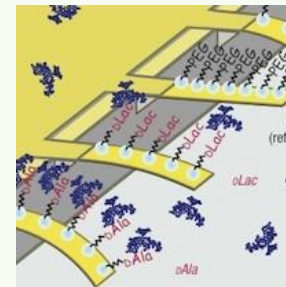
Biomaterials



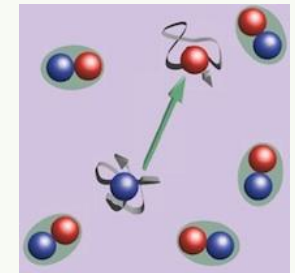
Quantum  
information



Optoelectronics



Sensors



Spintronics

# The future of nanoscience

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- ▶ The application of new extraordinary experimental tools has created an urgent **need for a quantitative understanding** of matter at nanoscale

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- › The **absence of quantitative models** that describe newly observed phenomena increasingly limits progress in the field

# The future of nanoscience

- » The application of new extraordinary experimental tools has created an urgent **need for a quantitative understanding** of matter at nanoscale
- » The **absence of quantitative models** that describe newly observed phenomena increasingly limits progress in the field
- » The absence of such tools would also **seriously inhibit wide-spread applications** ranging from molecular electronics to biomolecular materials

# The case of thermal conductivity

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$$J = -k \nabla T$$



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Is the Fourier's law applicable to nanostructures?

How size effects can affect thermal conductivity?

Can present theories interpret satisfactorily experiments?

# The case of thermal conductivity

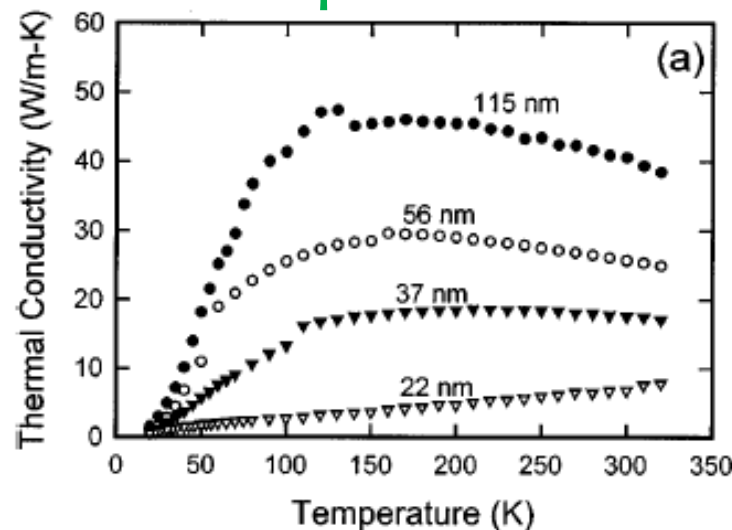
$$J = -k \nabla T$$

Is the Fourier's law applicable to nanostructures?

How size effects can affect thermal conductivity?

Can present theories interpret satisfactorily experiments?

An example: Si Nanowires



Measured thermal conductivity lower than bulk and diameter dependent

Quantum effects not negligible

Breakdown of Fourier's law

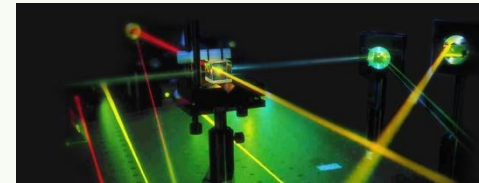
# Where do we need a theoretical effort?

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## Transport in nanostructures



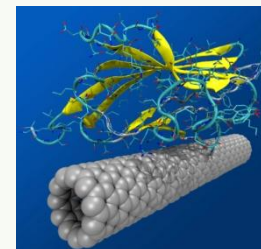
## Electronic and optical properties



## Coherence and decoherence tunneling



## Soft/hard matter interfaces



## Spintronics

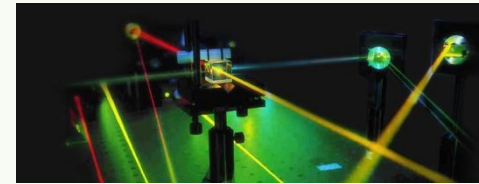


# Where do we need a theoretical effort?

## Transport in nanostructures



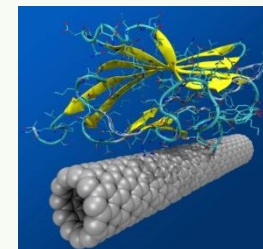
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# Materials design

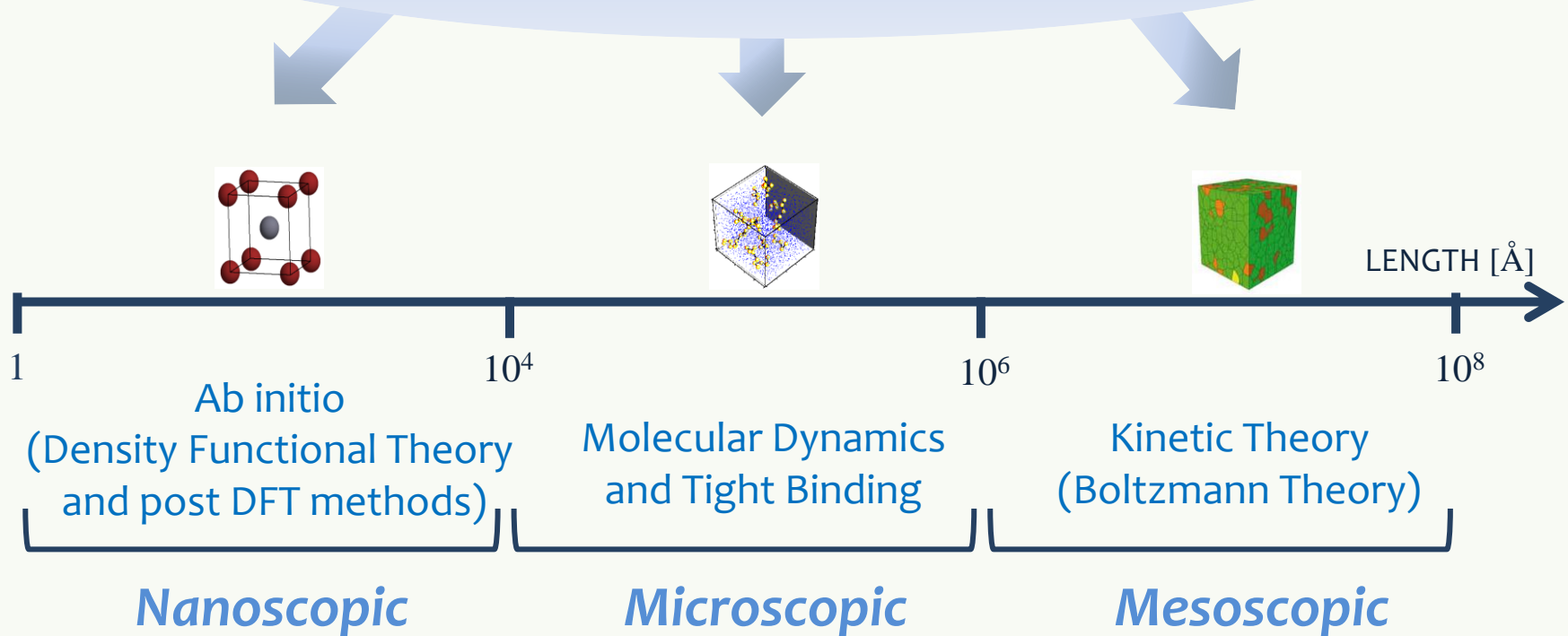
# Materials design

**Lengths scales of  
materials modelling**

Three blue arrows point downwards from the bottom of the central oval, indicating a flow or continuation of the topic.

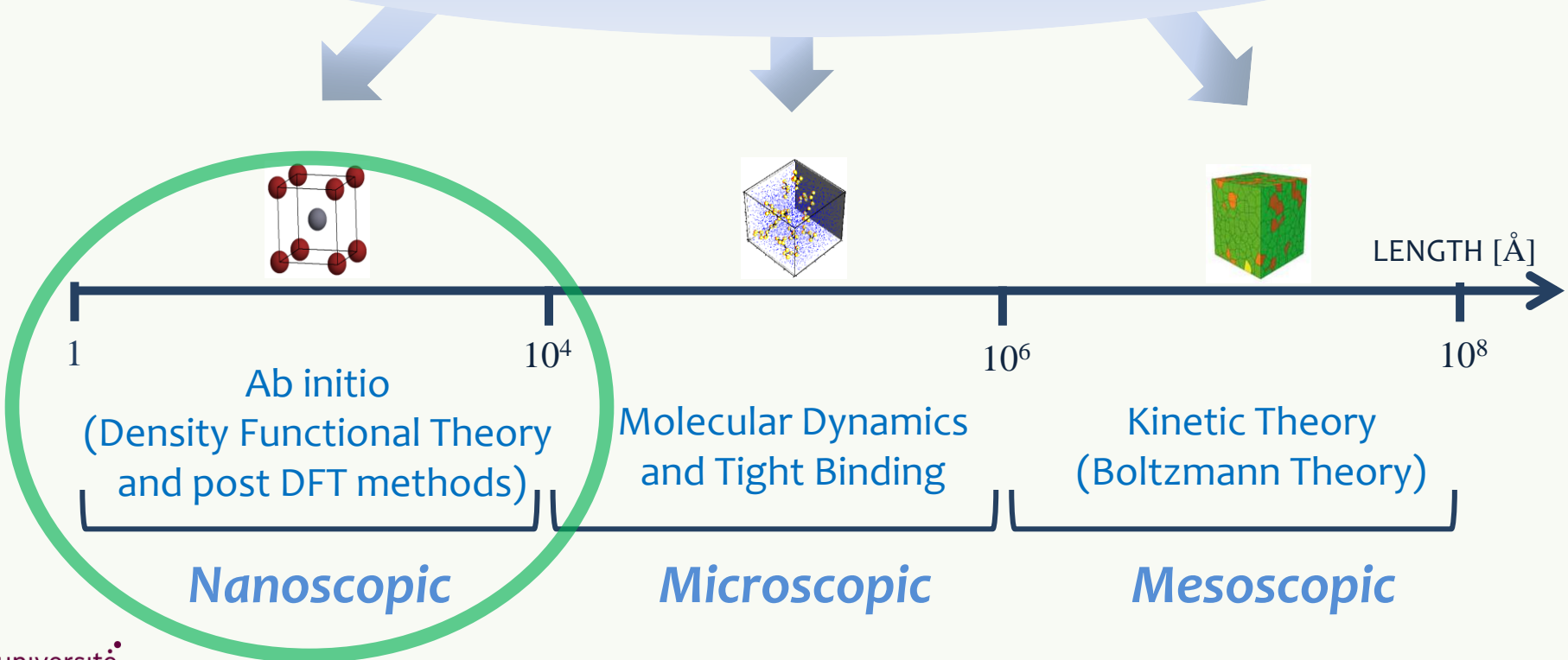
# Materials design

## Lengths scales of materials modelling



# Materials design

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# Ab initio computational modelling

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- ▶ The developement and use of mathematical models for describing and predicting certain properties of materials at a **quantitative level**
- ▶ ‘Ab initio’ or ‘from first principles’ refers to a bottom-up modelling strategy in which we **do not use any empirical parameters**
- ▶ Such kind of calculations are completely based on **quantum mechanics**, that can be considered as an engineering tool
- ▶ The complexity of such problems requires the use of **supercomputers**
- ▶ This is a discipline at the boundary between **materials science, physics and chemistry** on the one side, and **applied mathematics and software engineering**

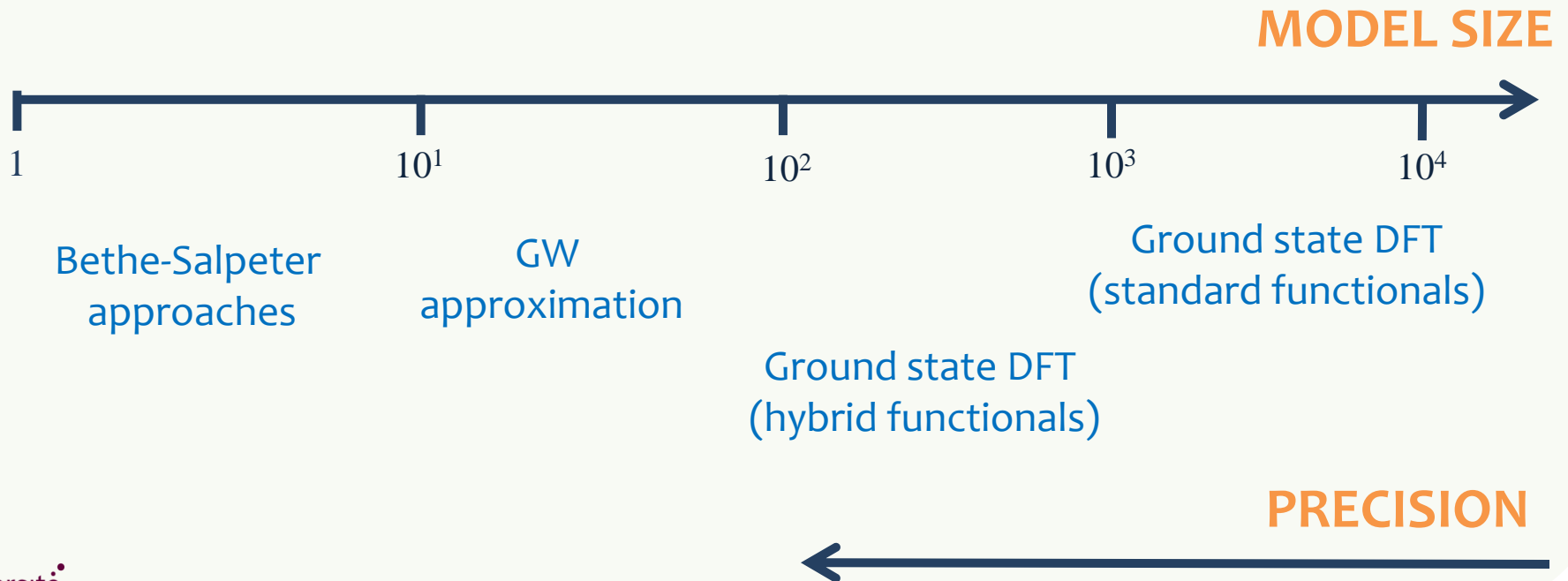
# Materials design

Materials design



# Materials design

## Precision of ab initio modelling



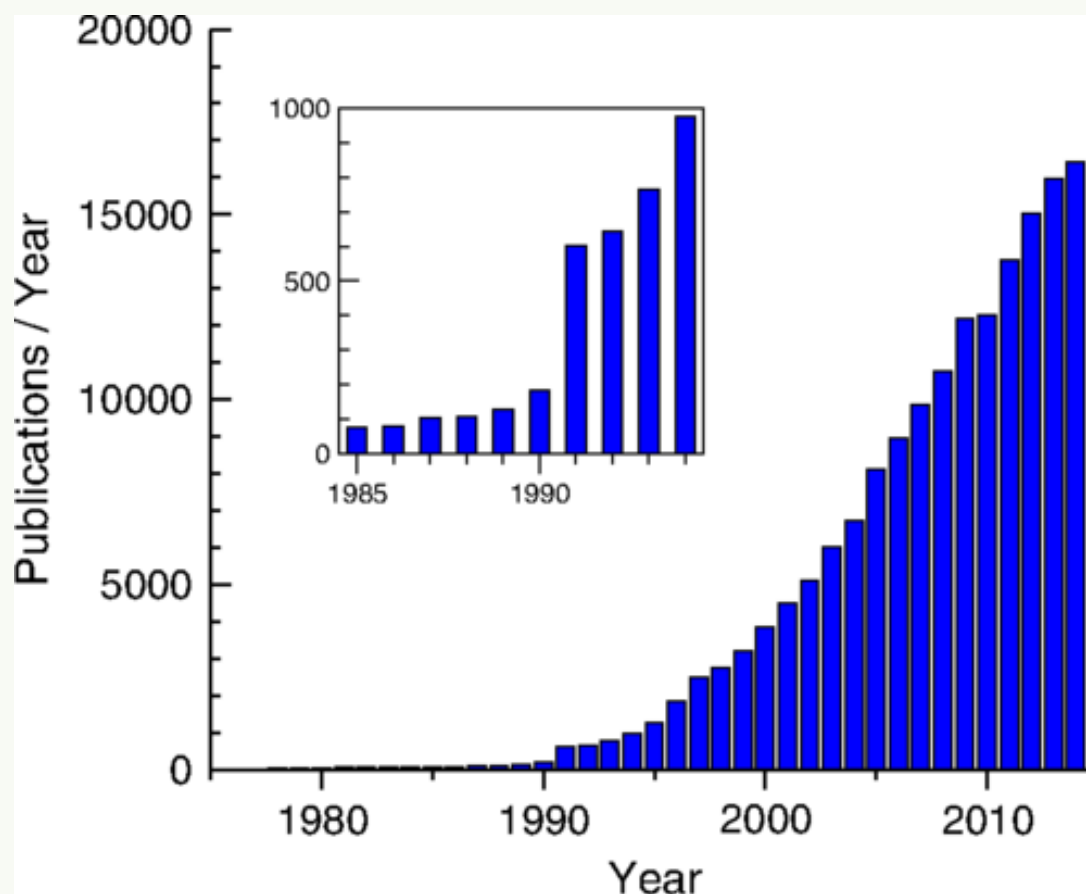
# Density Functional Theory

# Density Functional Theory

DFT is a very effective technique for studying molecules, nanostructures, solids, surfaces and interfaces by directly solving approximate versions of the Schrödinger equation

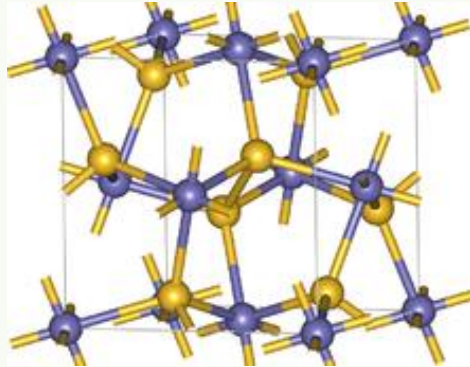
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# The many body electronic structure

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$N_e$  electrons  
 $N_n$  nuclei

Schrödinger equation for interacting particles

$$\hat{H}\Psi(\{R\}, \{r\}) = E\Psi(\{R\}, \{r\})$$

$$\begin{aligned}\hat{H} = & \hat{T}_n(\{R\}) + \hat{V}_{nn}(\{R\}) + \hat{T}_e(\{r\}) + \\ & + \hat{V}_{ee}(\{r\}) + \hat{U}_{en}(\{R\}, \{r\})\end{aligned}$$

# The many body hamiltonian

# The many body hamiltonian

$$\hat{T}_n = \sum_{I=1}^{N_n} -\frac{\nabla_I^2}{2M_I}$$

Kinetic energy of nuclei

$$\hat{T}_e = \sum_{i=1}^{N_e} -\frac{\nabla_i^2}{2m_e}$$

Kinetic energy of electrons

$$\hat{V}_{nn} = \frac{1}{2} \sum_{I,J,I \neq J}^{N_n} \frac{Z_I Z_J}{|R_I - R_J|}$$

Ion-ion interactions

$$\hat{V}_{ee} = \frac{1}{2} \sum_{i,j,i \neq j}^{N_e} \frac{1}{|r_i - r_j|}$$

Electron-electron interactions

$$\hat{U}_{en} = - \sum_{j,J}^{N_e, N_n} \frac{Z_J}{|R_J - r_j|}$$

Electron-ion interactions

Exactly solvable for two particles (analytically)  
and very few particles (numerically)

How to deal with  $N \sim 10^{23}$  particles?

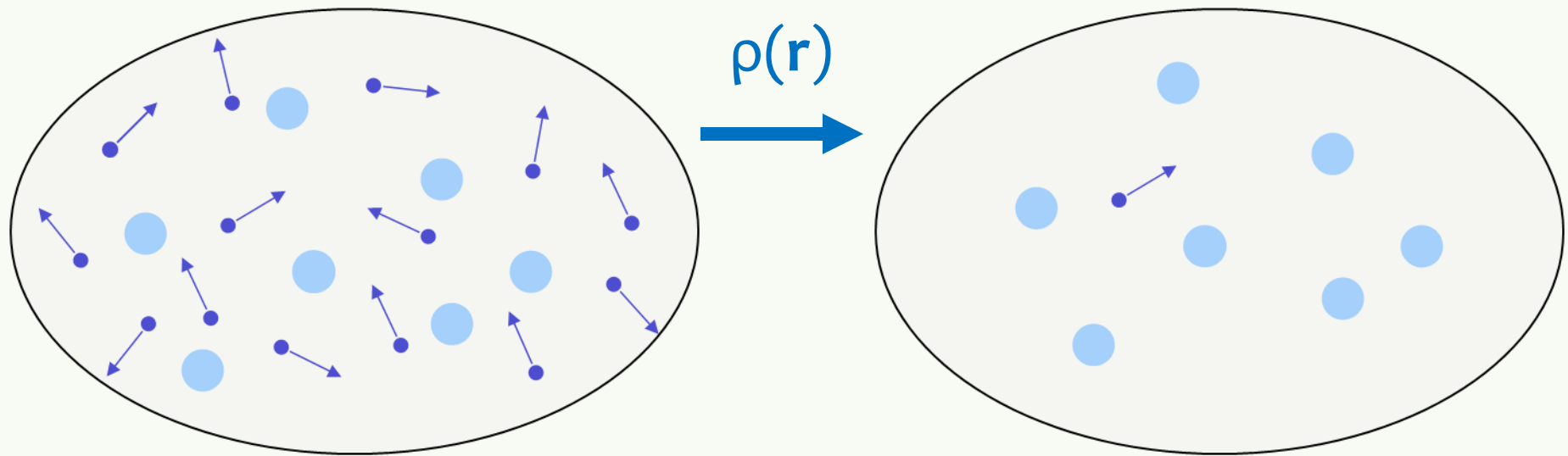


# Density Functional Theory (DFT)

# Density Functional Theory (DFT)

It can map, exactly, the **interacting** problem to a **non-interacting** one

Hohenberg and Kohn, Phys. Rev. B 136, 864 (1964)



interacting particles in a real  
external potential

a set of non-interacting electrons (with  
the same density as the interacting  
system) in some effective potential

# Popularity of DFT

# Popularity of DFT

- ▶ **Transferability:** same technique for describing different classes of materials
- ▶ **Simplicity:** based on simply and intuitive equations
- ▶ **Realiability:** possibility of making direct and quantitative comparison with experiments
- ▶ **Software sharing:** online platforms and adoption of open-source software model
- ▶ **Reasonable starting point:** even when it fails in describing correctly a property, it respresent an accurate starting point for more accurate theory (more computationally demanding)

# Ab initio properties of materials

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## ELECTRONIC STRUCTURE

Density Functional Theory (DFT),  
Band structure, Density of states,  
Wave function

## CHEMICAL REACTIVITY

Nudge Elastic Bands  
Activation barrier,  
adsorption energies

## OPTICAL RESPONSE

Many Body Perturbation Theory,  
Optical Absorption, Excitons,  
GW approach, EELS

## VIBRATIONAL PROPERTIES

Density Functional  
Perturbation Theory (DFPT),  
Raman spectra, phonons, phase diagrams

## ELASTIC PROPERTIES

DFT  
Elastic constants and  
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# Theoretical method and codes



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**SIESTA code** <http://icmab.es/siesta>

## Density Functional Theory (DFT)

- Exchange-correlation functional: LDA
- Optimized double-z polarized basis-set
- Troullier-Martins pseudopotentials



**VASP code** <http://www.vasp.at>

## Hybrid - DFT

- Exchange-correlation functional: PBE-HSE06
- Plane-wave basis set
- PAW approach pseudopotentials



**YAMBO code** [www.yambo-code.org](http://www.yambo-code.org)

## GW and optical properties calculations

- GW one shot perturbative approach
- Plane-wave basis set
- RPA and BSE approaches

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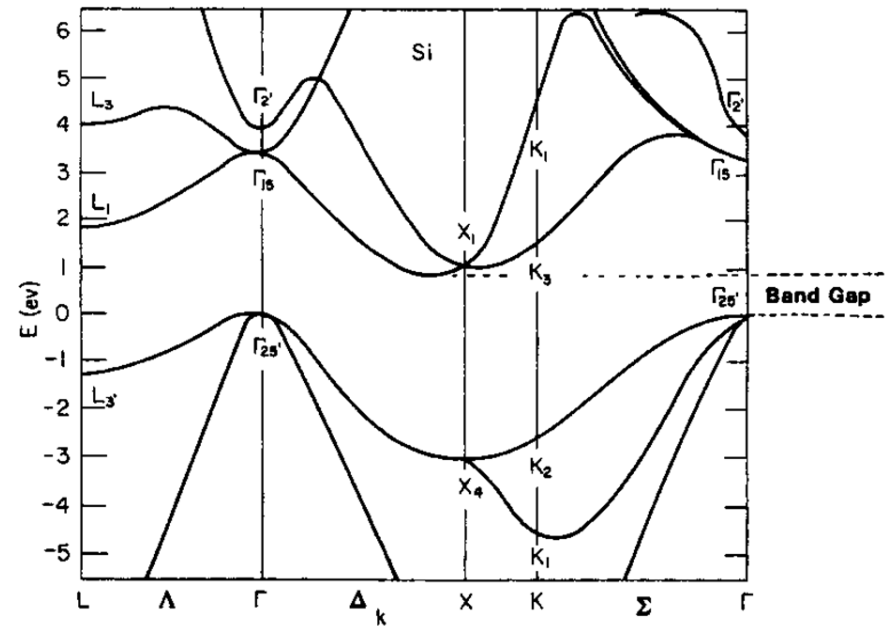
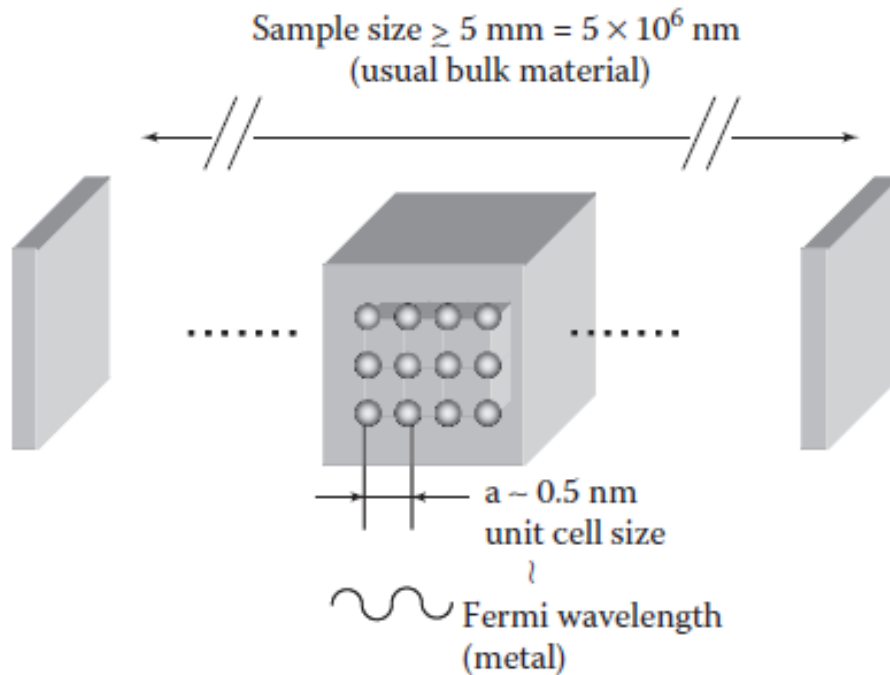
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# Space for electrons in materials

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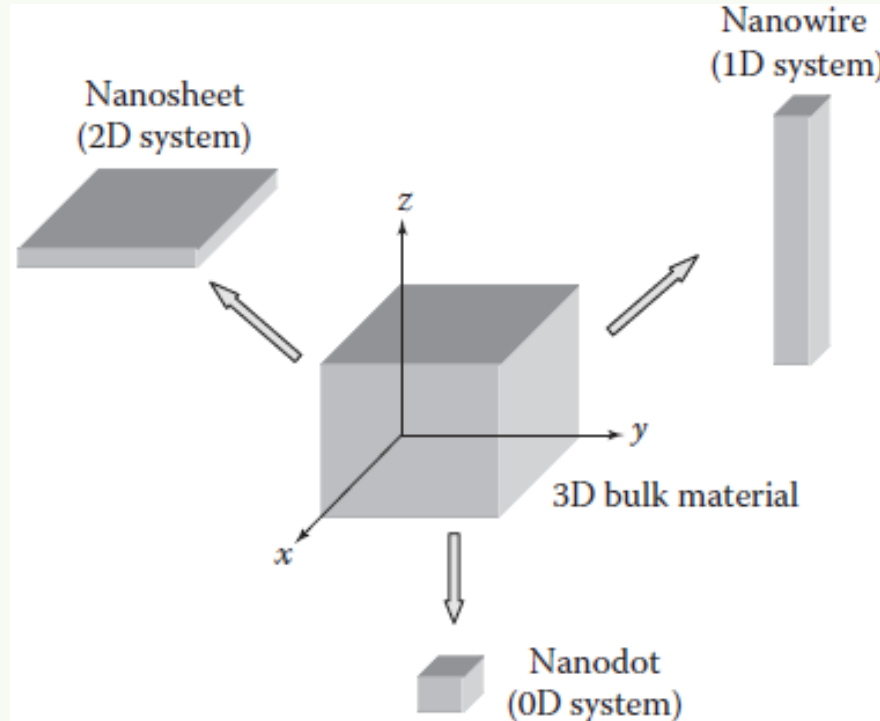
Si band structure

T. Tsurumi et al. *Nanoscale physics for materials science*, Taylor and Francis (2010)

- In macroscopic materials the Fermi wave length is of the same order as the unit cells
- Thus the inner space is too large for electrons to feel boundaries
- The possible energy levels for electrons  $E(k) = (\hbar^2/2m^*)k^2$  are dense with very small energy spacing

# Quantum confinement effect (QCE)

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T. Tsurumi et al. *Nanoscale physics for materials science*, Taylor and Francis (2010)

Quantum confinement effect (QCE) is defined as: a reduction in the degrees of freedom of the carrier particles, implying a reduction in the allowed phase space

E.G. Barbagiovanni et al. *Appl. Phys. Rev.* 1, 011302 (2014)



# Consequences of QCE

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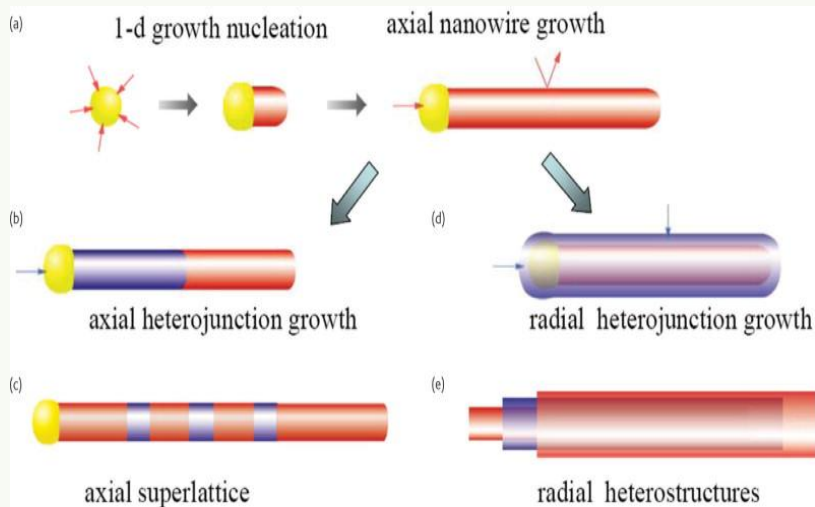
- Electrons traveling along the direction that has been reduced in size reach the boundaries and are confined
- This confinement causes the quantization of the electron wavelength and of the energy spectrum
- The contribution of edge-localized surface states can become relevant
- Size effects dominate the physics of nanostructures

# The case of Si nanowires (Si NWs)

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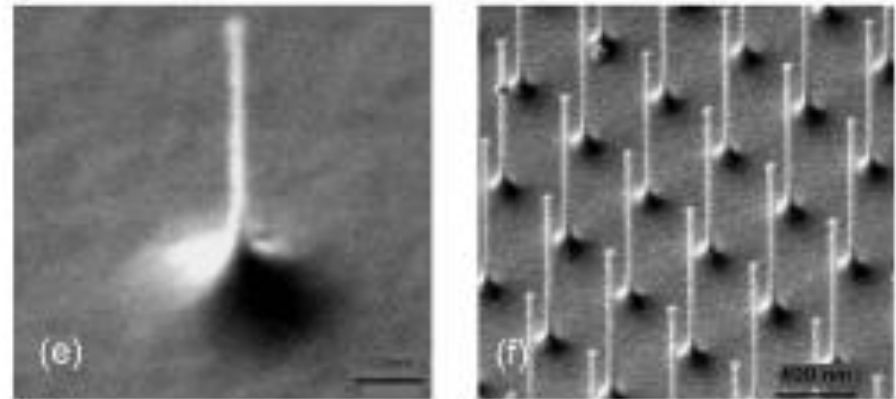
One-dimensional nanostructures with precise composition, morphology, interfaces and electrical properties ( $d=2-3$  nm)

## Vapor-Liquid-Solid (VLS) growth



O. Hayden et al., Nanotoday 3 5-6 (2008)

## Lithography and etching processes

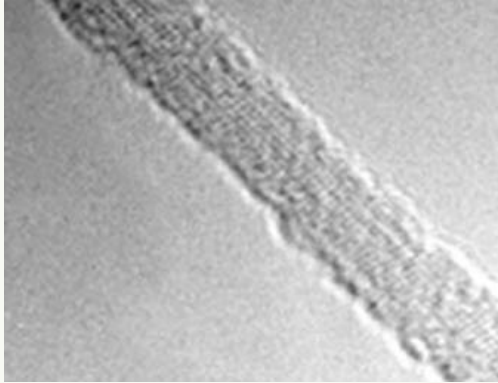


N. Singh et al., IEEE Trans. Electron Devices 55, 11 (2008)

# QCE in Si NWs: electronic structure

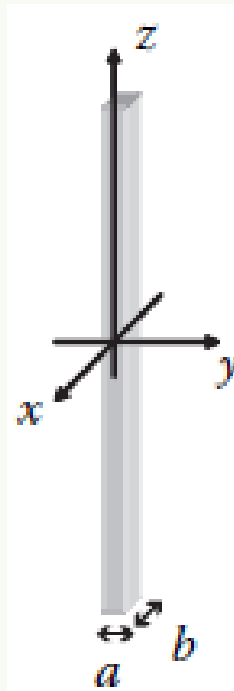
# QCE in Si NWs: electronic structure

## Real system



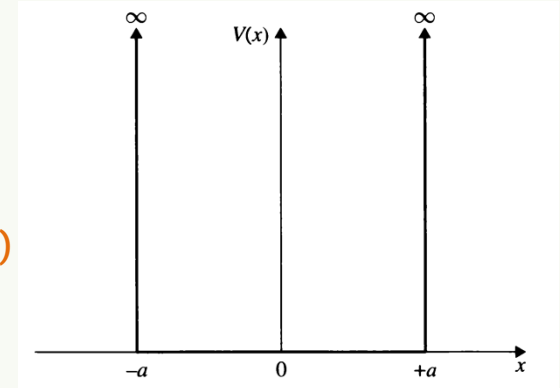
Wu et al., Nano Lett. 4, 433 (2004)

## Geometrical model



T. Tsurumi et al., Taylor and Francis (2010)

## Infinite potential well (particle-in-a-box model)



B.H. Branden and Joachin,  
*Quantum Mechanics* (2000)

# QCE in Si NWs: electronic structure

# QCE in Si NWs: electronic structure

The motion of electrons is restricted to be in the direction of confinement. Their kinetic energy increases and the eigenstates energies are given by:

$$E_n = \frac{\hbar^2 n^2 \pi^2}{2m^* d^2}$$

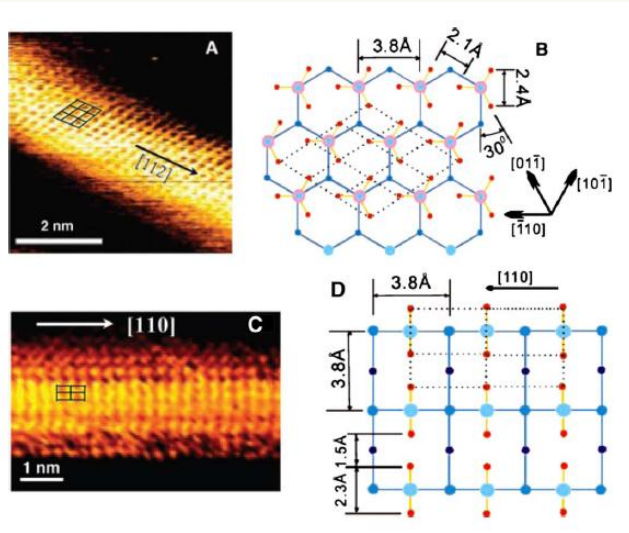
m\*=effective mass  
d=width of the well

Not only the energy levels but even the spacing between them increase when the diameter is reduced (QCE increases)

QCE has a dramatic effect on semiconductor NWs (like Si NWs) because it affects the **energy band gap**

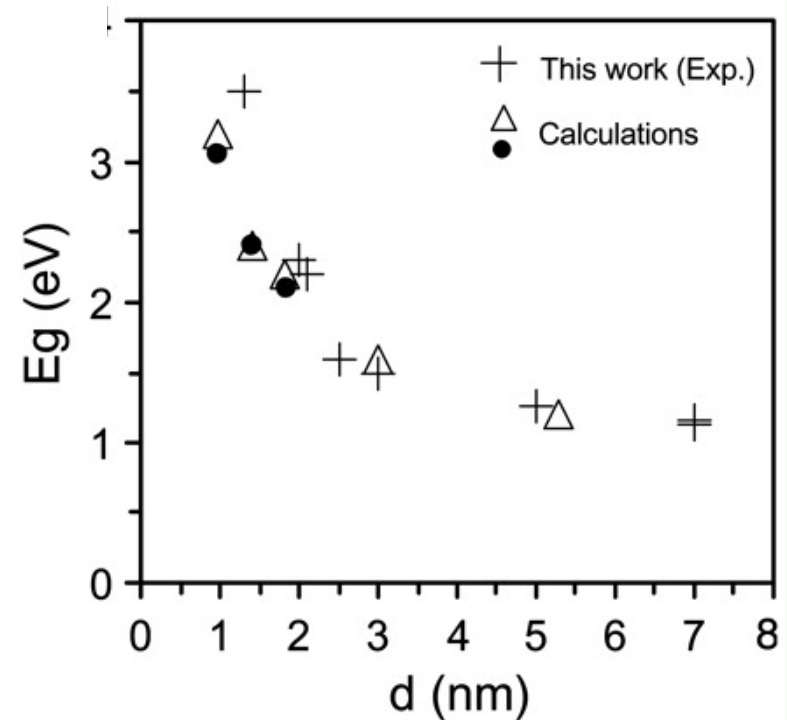


# QCE in Si NWs: electronic structure



Ma, D. D. D. et al., Science 299, 1874 (2003)

## STS measurements



Ma, D. D. D. et al., Science 299, 1874 (2003)

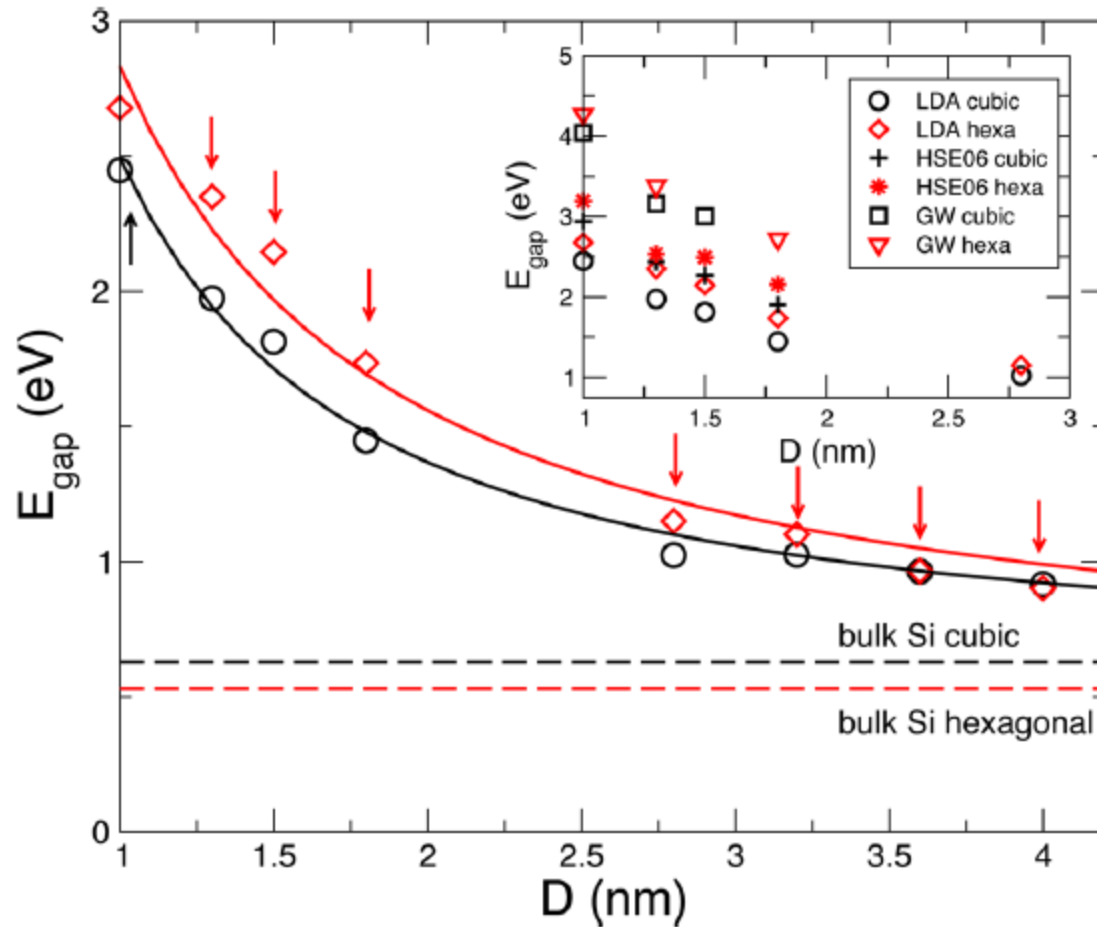
$$E_{\text{gap}} = E_{\text{bulk, gap}} + C(1/d)^{\alpha}$$

with  $\alpha=0.9 - 1.1$

M. Bruno et al., Phys. Rev. Lett. 98, 036807 (2007)

# QCE in Si NWs: DFT scaling of the gap

# QCE in Si NWs: DFT scaling of the gap



DFT – LDA

$$\alpha_{\text{hex}} = 1.16$$

$$\alpha_{\text{cub}} = 1.34$$

$$E_{\text{gap}} = E_{\text{bulk, gap}} + C(1/d)^{\alpha}$$

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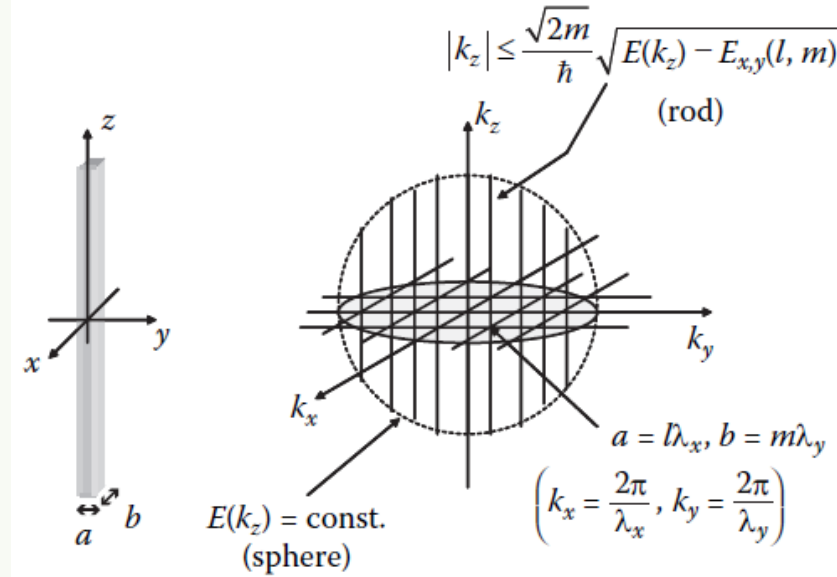
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# QCE in NWs: energy quantization

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T. Tsurumi et al., Taylor and Francis (2010)

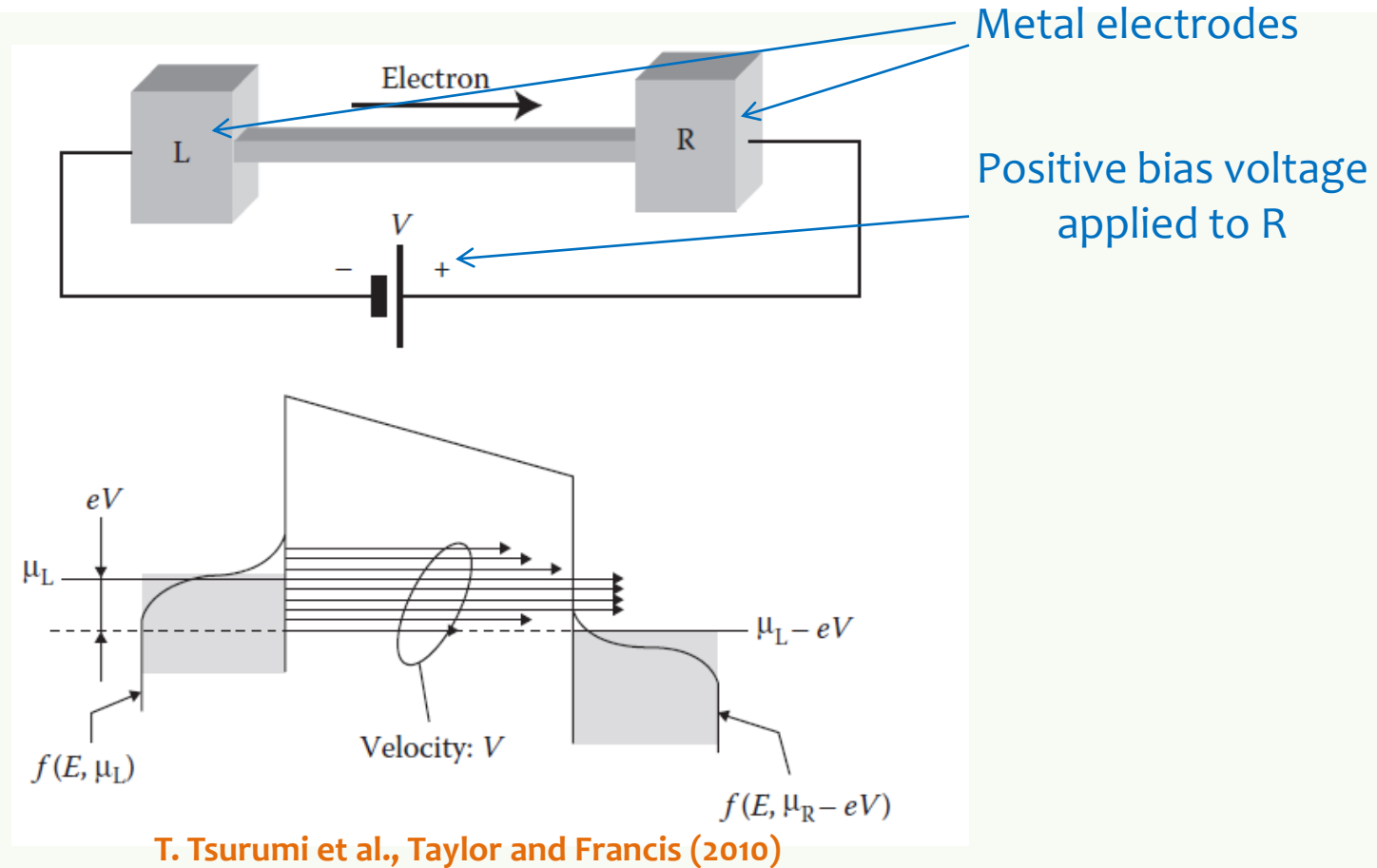
The total energy of electrons is given by the formula:

$$E(k_z) = \underbrace{\frac{\hbar^2}{2m^*} k_z^2}_{\text{Kinetic energy}} + \underbrace{E_{x,y}(l, m)}_{\text{Quantized energy in the plane}}$$

# QCE in NWs: Landauer approach



# QCE in NWs: Landauer approach

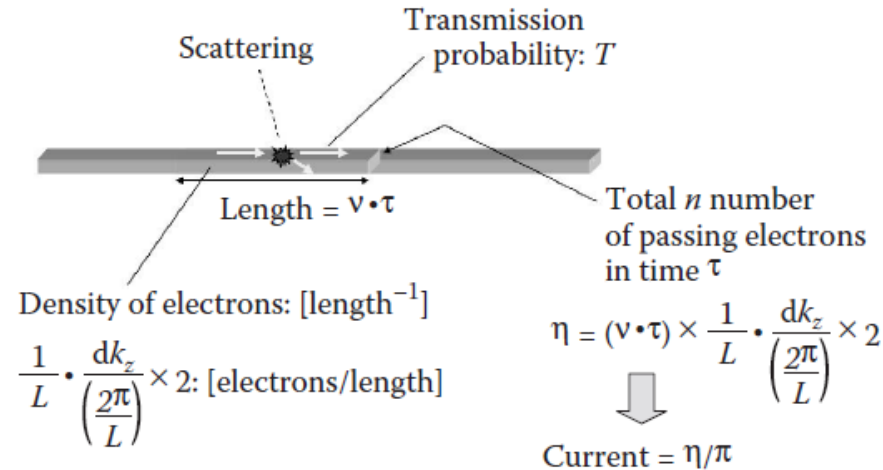


Electrons move with a group velocity  $v(k)$  and transmission probability  $T(k)$

Can we calculate the current between electrodes?

# QCE in NWs: quantized conductance

# QCE in NWs: quantized conductance



T. Tsurumi et al., Taylor and Francis (2010)

## Current

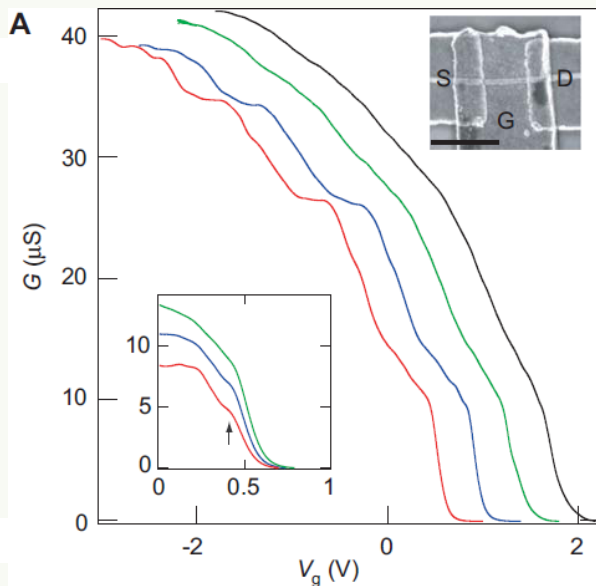
$$I = \frac{2e}{h} \int_0^\infty \{f(E, \mu) - f(E, \mu - eV)\} T(E) dE$$

## Conductance (if the bias is small)

$$G = \frac{2e^2}{h} \int_0^\infty \left\{ -\frac{\partial f(E, \mu)}{\partial E} \right\} T(E) dE \equiv \frac{2e^2}{h} T_\mu$$

# QCE in NWs: quantized conductance

- In 1D systems the conductance of electrons is quantized in units of  $2e^2/h$
- Electrons states along the wire are associated with quantized states in the plane
- Each of the quantum states in the plane has equal unit of conductance  $(2e^2/h)T_\mu$  along the wire axis
- The conductance of NWs is hence the product of the number of quantum states and their quantized conductance  $(2e^2/h)T_\mu$

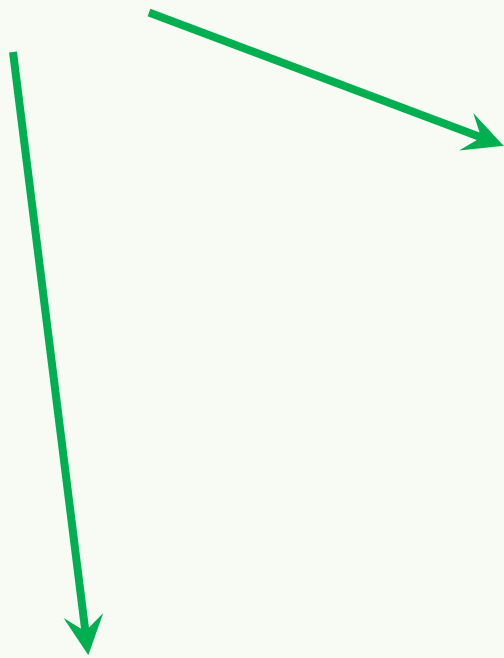


**Experimental  
demonstration on  
SiGe core-shell NWs**

# QCE in Si NWs: DFT+Landauer approach

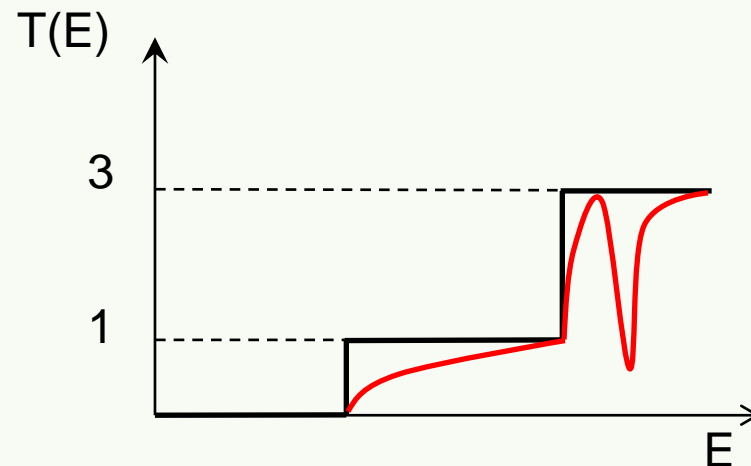
# QCE in Si NWs: DFT+Landauer approach

Conductance is calculated in terms of transmission probability  $T(E)$  through the available transmitting channels



For a pristine Si NW:  $T(E) = 1$   
For a defected NW:  $T(E) < 1$

How many?  
For an infinitely long NW: as many as electron states at that energy



M. Amato et al. Nano Lett. 12, 2717 (2012)

# Outline

- › Introduction on nanoscience
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- › Quantum confinement in Si nanowires (Si NWs)
  - Electronic structure
  - Transport properties
  - Doping effects
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  - Crystal phase effects on bulk
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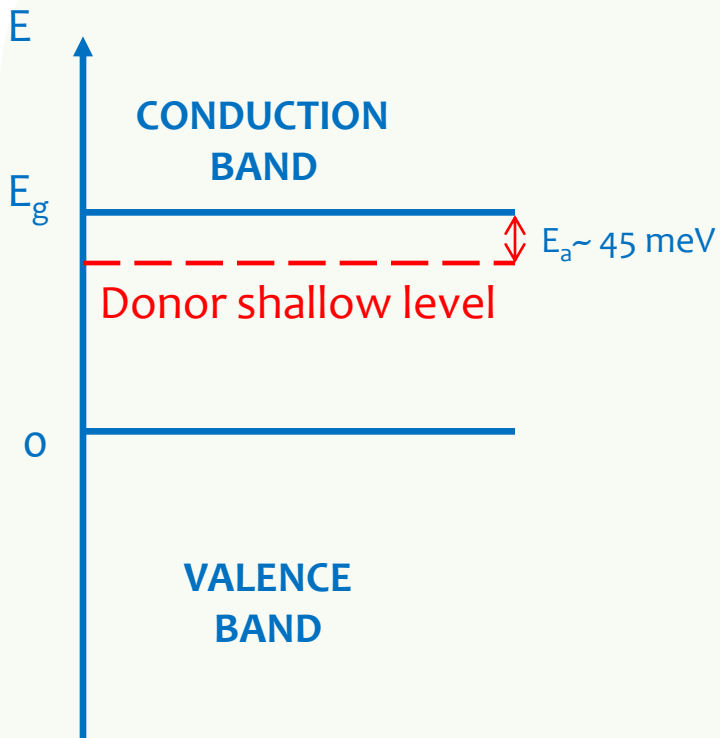
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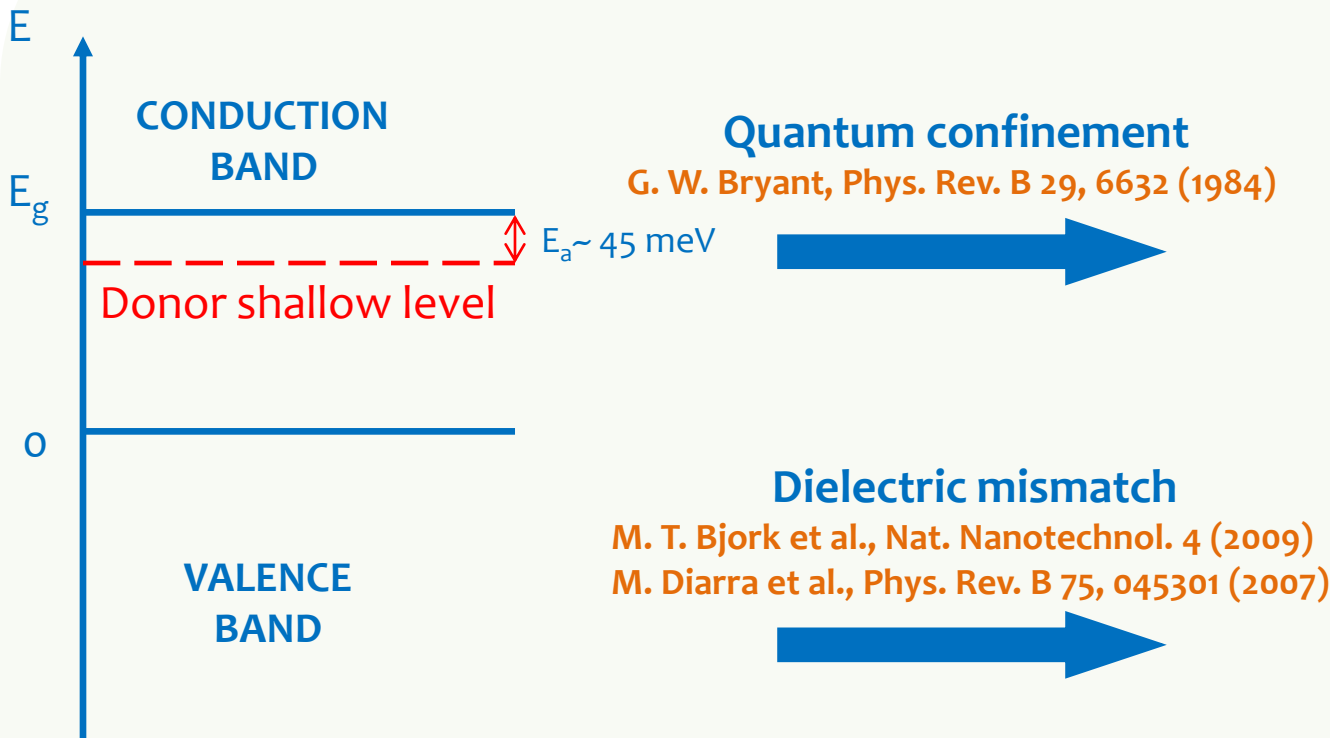
Si bulk P-doped



- Ionization energy  $\sim$  few hundredths of eV
- Ionized impurity at room temperature
- High doping efficiency

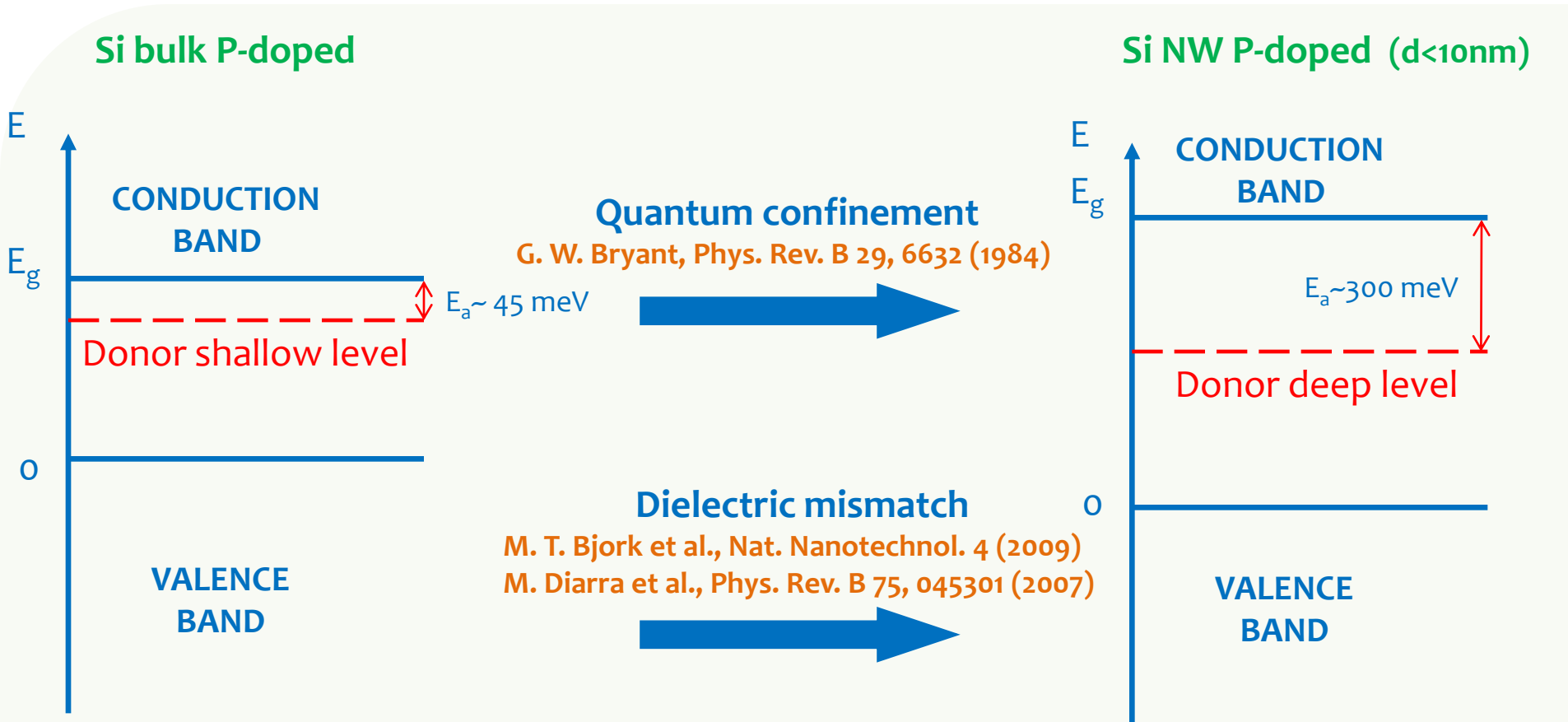
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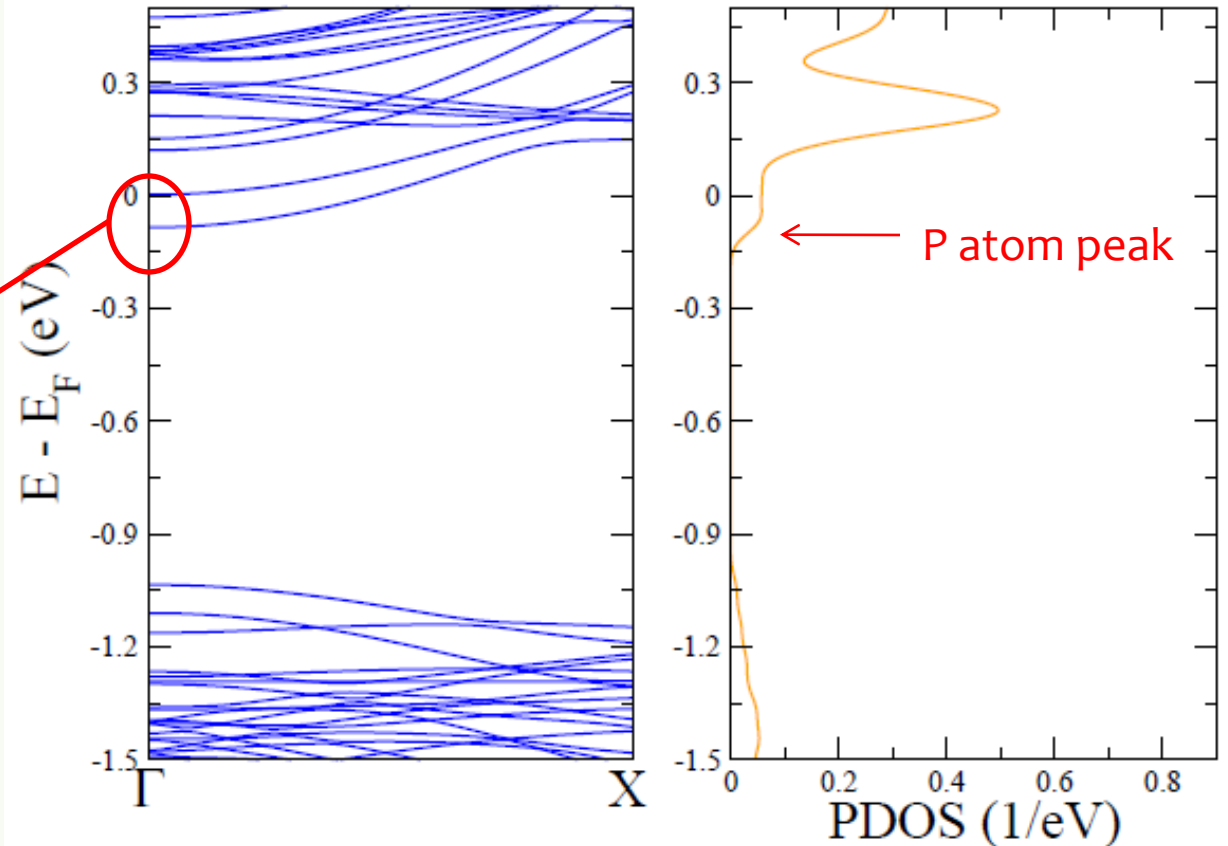
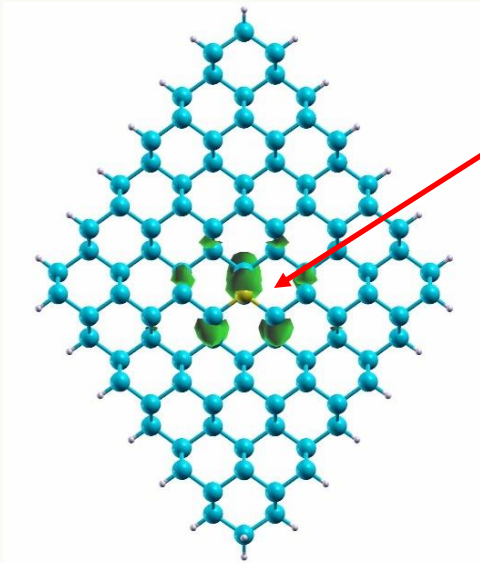


- Ionization energy  $\sim$  few hundredths of eV
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- Ionization energy increase
- Donor level deep into the band gap
- Doping deactivation with respect to bulk

# QCE in Si NWs: P deactivation in DFT

# QCE in Si NWs: P deactivation in DFT



M. Amato et al., Nano Lett. 11, 594–598 (2011)

P-impurity state deep into the band gap  
Activation energy is too high ( $\sim 100$  meV)

# Some comments on QCE in Si NWs



# Some comments on QCE in Si NWs

The diameter is a powerful tool to modulate physical properties

- **Electronic Structure.** Reducing the size can induce an opening of the bulk band gap that can cause a modulation of optical properties
- **Transport.** The quantization due to the confinement leads to a quantization of the electrical conductance. Si NWs are hence quantum confined ballistic conductors
- **Doping.** Intentional addition of impurities in Si NWs is not efficient as for bulk systems due to the impurity deactivation

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Is it possible to further modify  
the physical properties of Si NWs?

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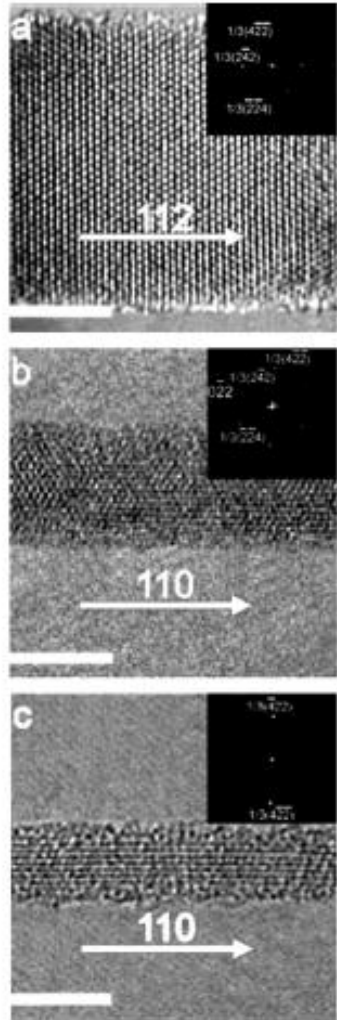
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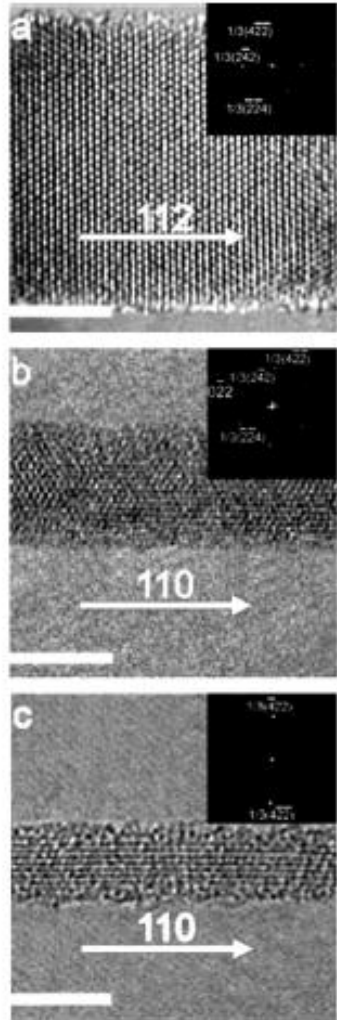
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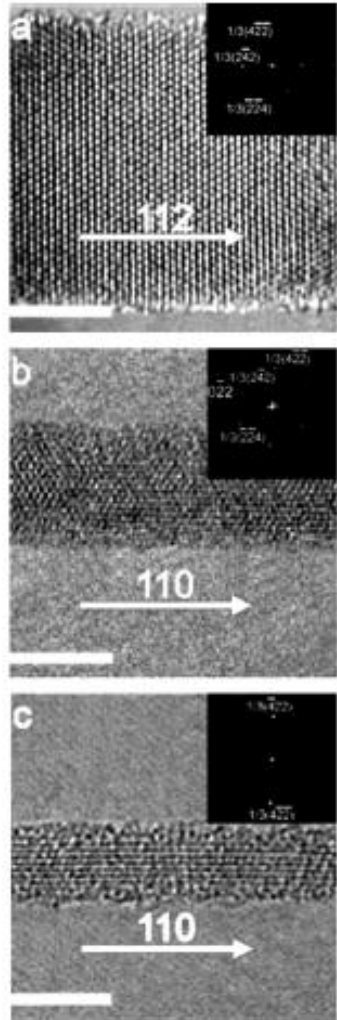
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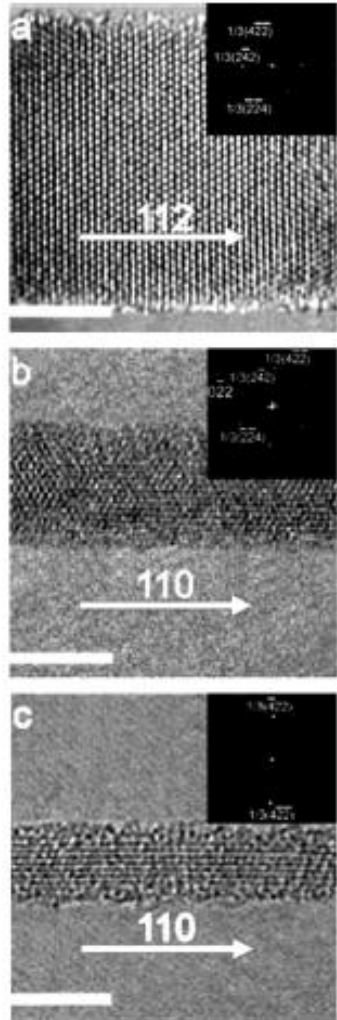
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# The crystal phase of group IV NWs



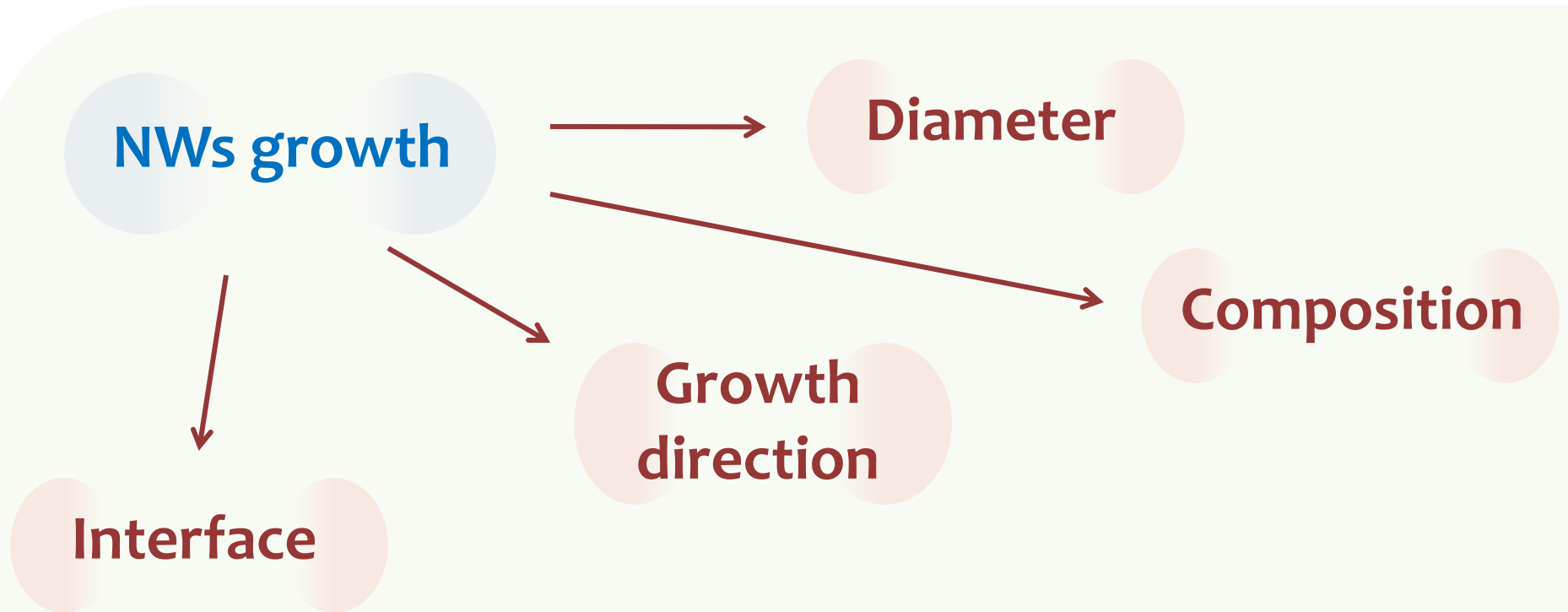
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- ▶ In particular it was found that they are rodlike structures constructed around a bulk Si single crystalline core
- ▶ This finding was mostly independent of the growth method, although it has been demonstrated that the quality of the crystal and the density of defects can depend on it
- ▶ All these experiments confirmed that, as can be intuitively expected, **Si NWs crystallize in the diamond-like cubic structure as in their bulk counterpart (3C phase)**

# Properties of group IV NWs

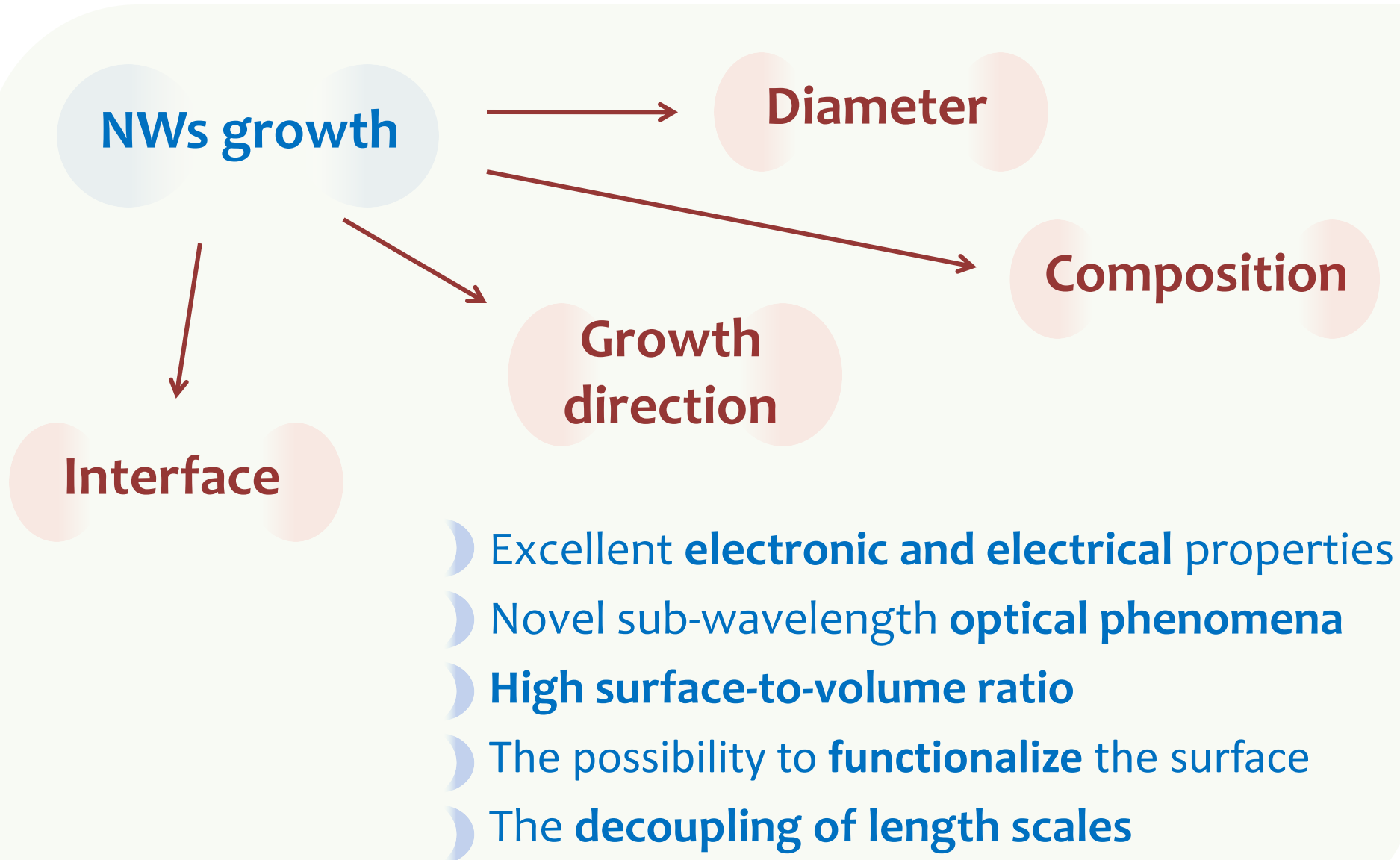
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NWs growth

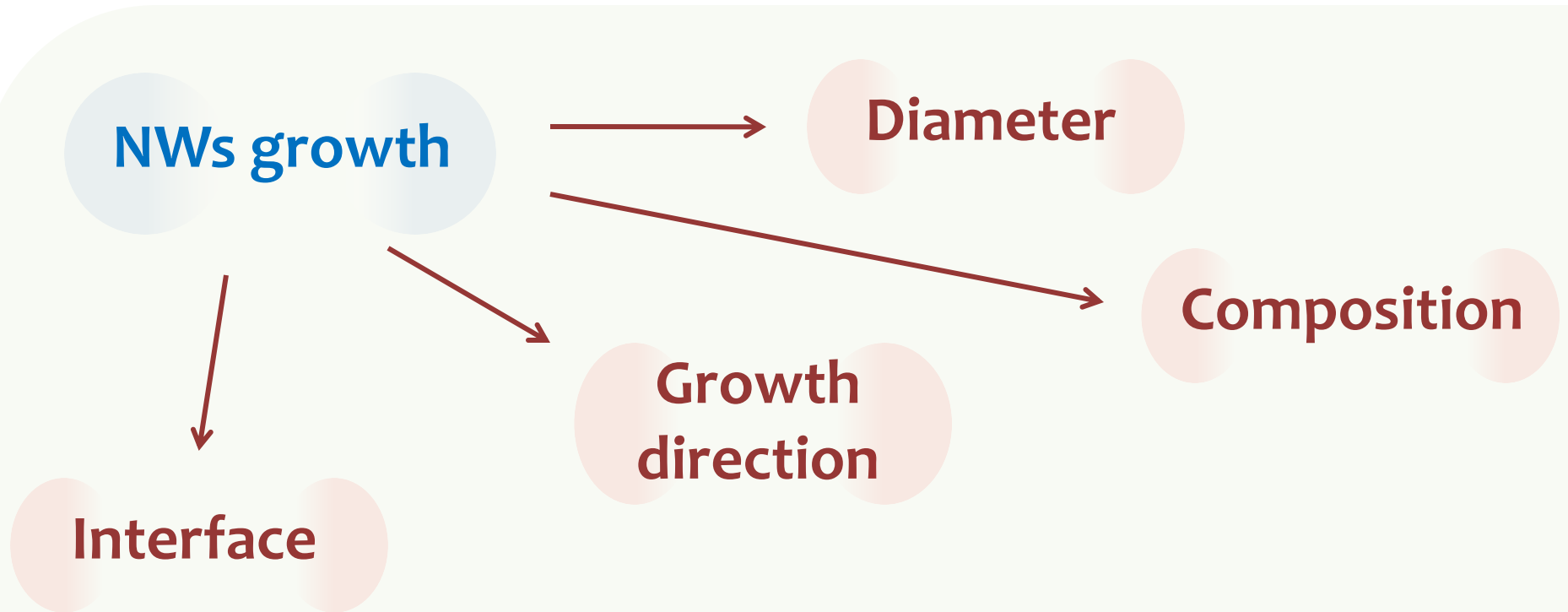
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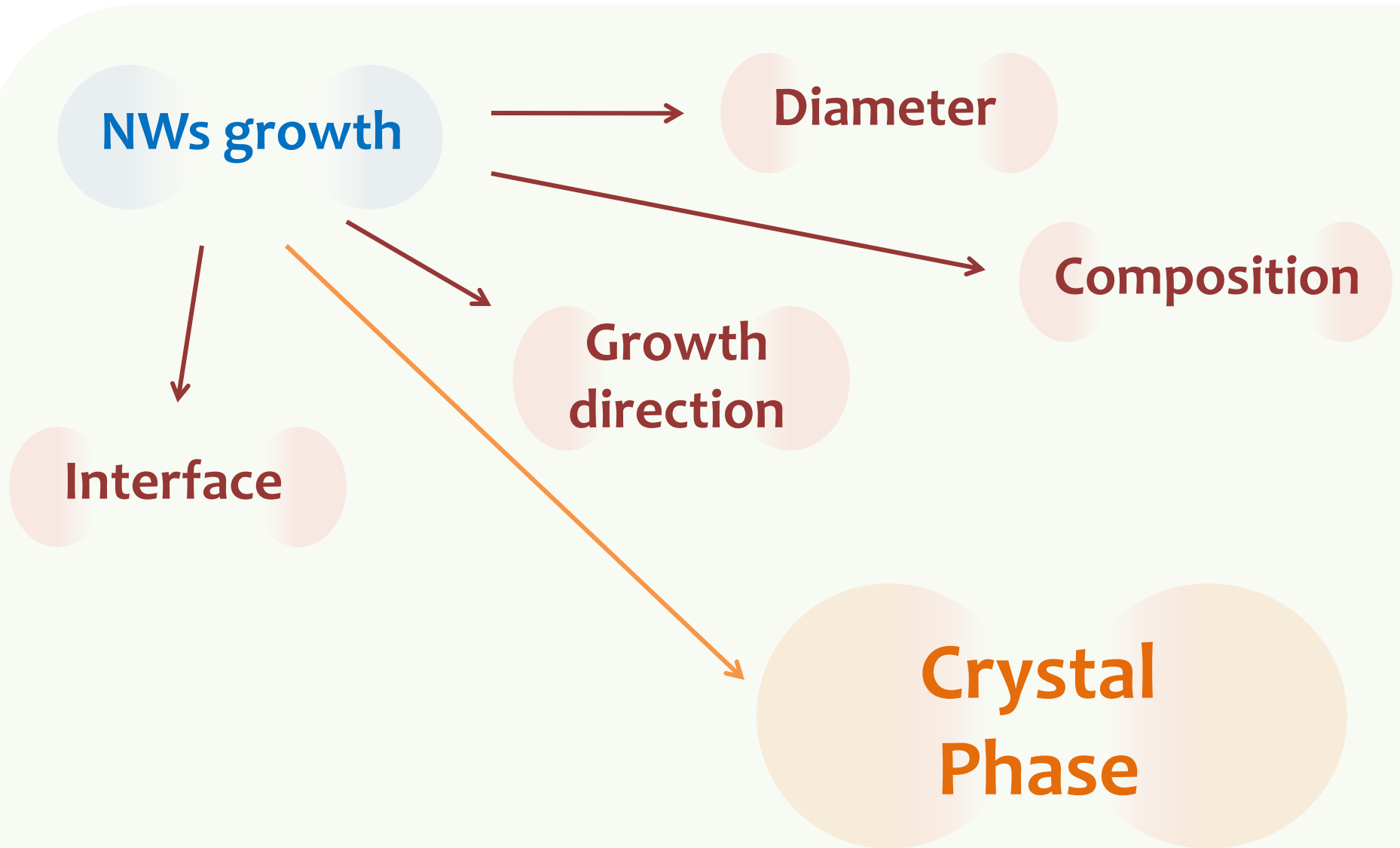
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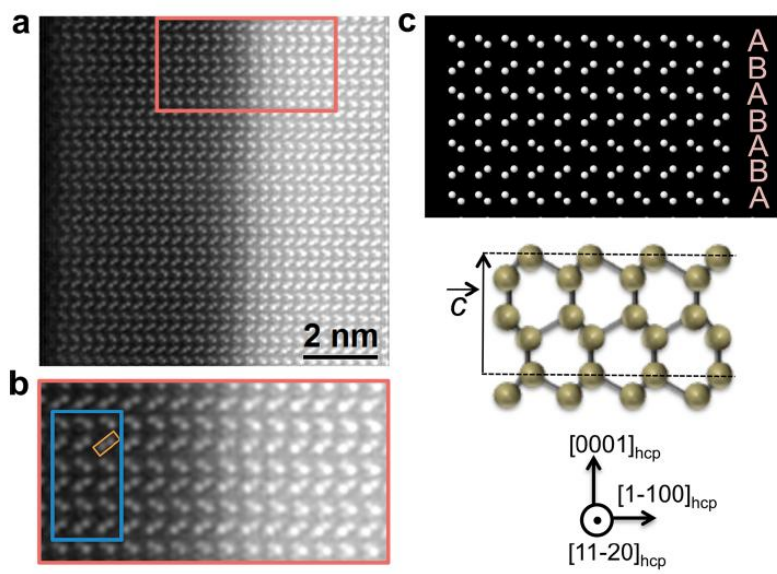


# Hexagonal-diamond (2H) NWs



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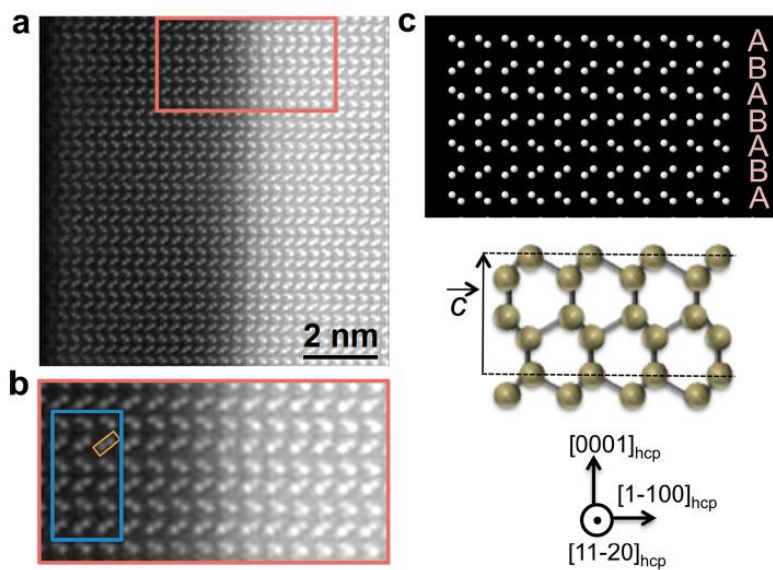
## Hexagonal Si nanoshells on top of hexagonal GaP



H. T. Hauge et al., Nano Lett. 15, 5855 (2015)

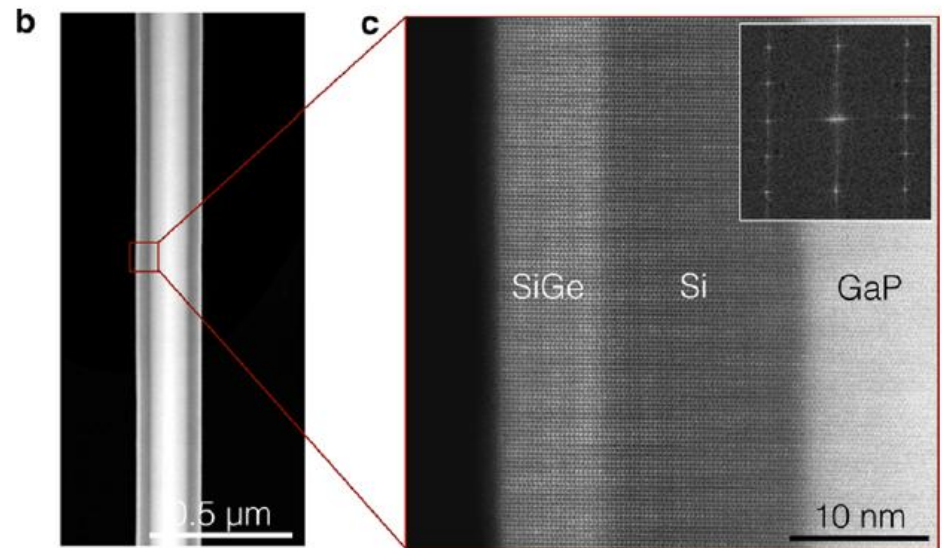
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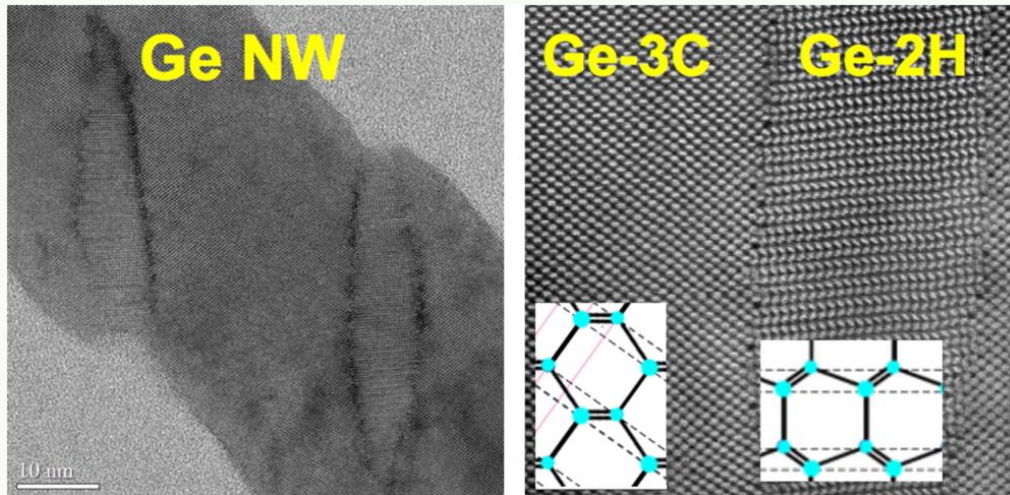


H. T. Hauge et al., Nano Lett. 17, 85-90 (2017)

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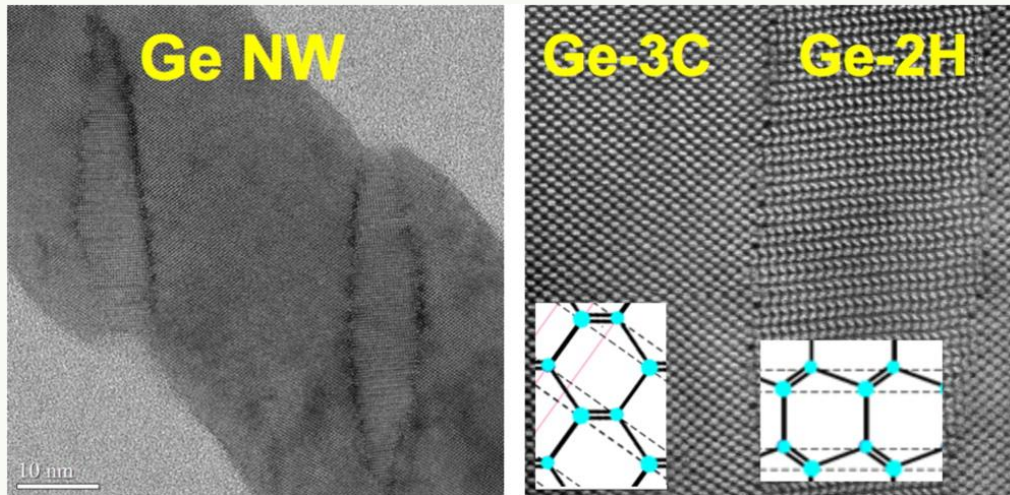
## Hexagonal-cubic diamond Ge NWs homojunctions



L. Vincent et al., Nano Lett. 14, 4828 (2014)

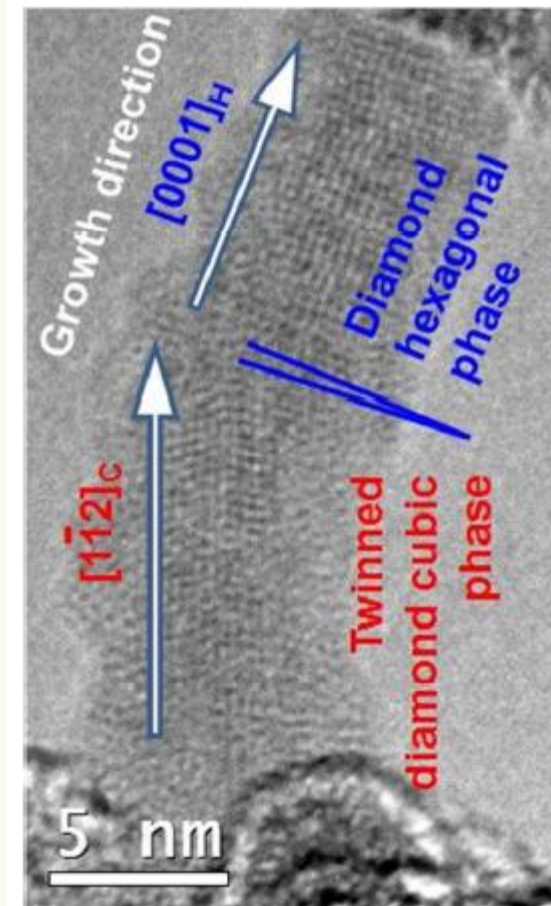
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L. Vincent et al., Nano Lett. 14, 4828 (2014)

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J. Tang et al., Nanoscale 9, 8113 (2017)

# A not obvious result for bulk



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Bulk hexagonal diamond Si metastable  
and obtained only at high pressures



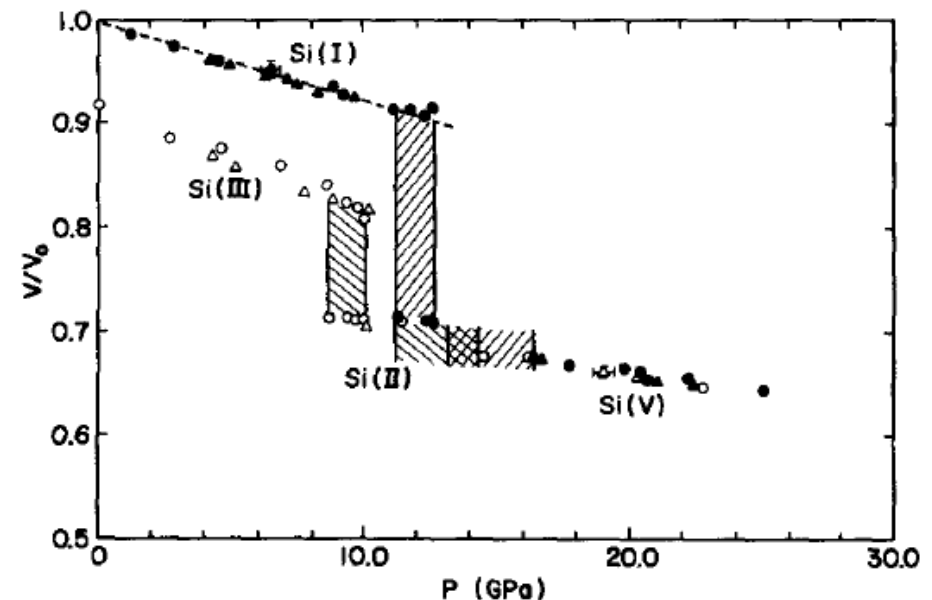
Pirouz et al., Acta Metall. Mater. 38, 313 (1990)

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Pirouz et al., Acta Metall. Mater. 38, 313 (1990)



J. Z. Hu and I. L. Spain, Sol. State Comm. 51, 263 (1984)



# Why hexagonal NWs are interesting?

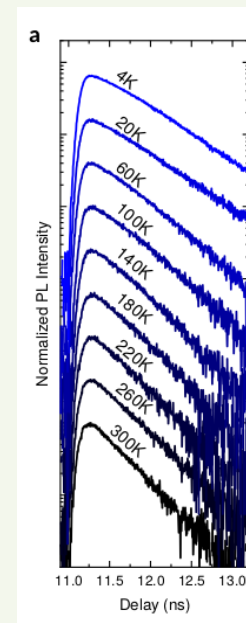
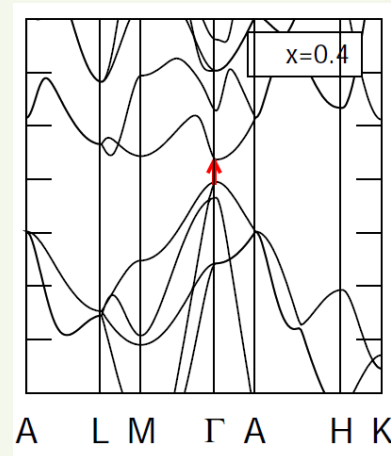
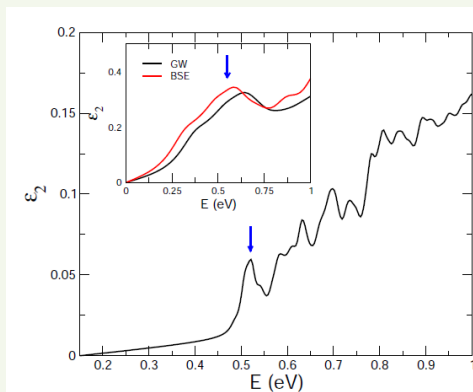
# Why hexagonal NWs are interesting?

- › Modified **band gap**, **band splittings** and **effective masses**
- › Modulation of the **optical absorption** spectra
- › Greater **electronic mobility**
- › Possibility to modulate **band offset** type and magnitudes
- › Control of **electron-hole** separations

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## Example. 2H-SiGe NWs as light emitters



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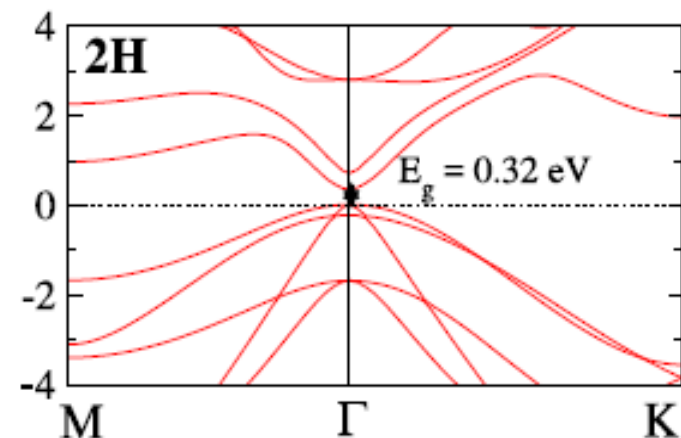
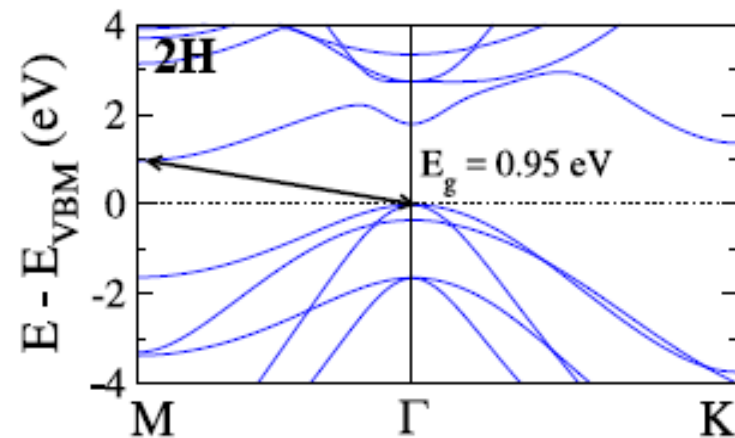
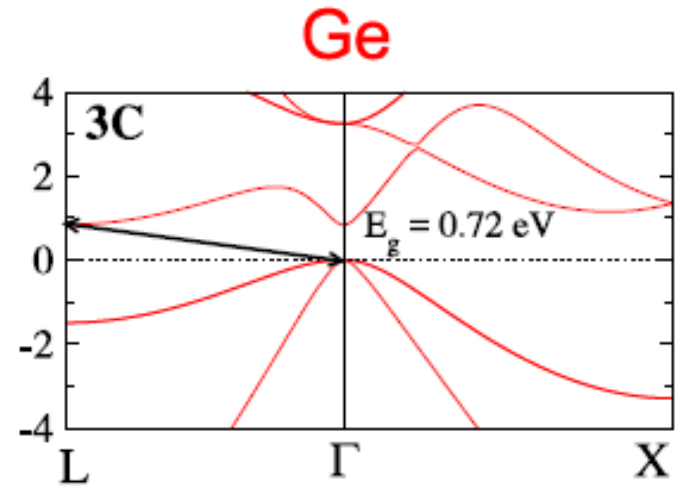
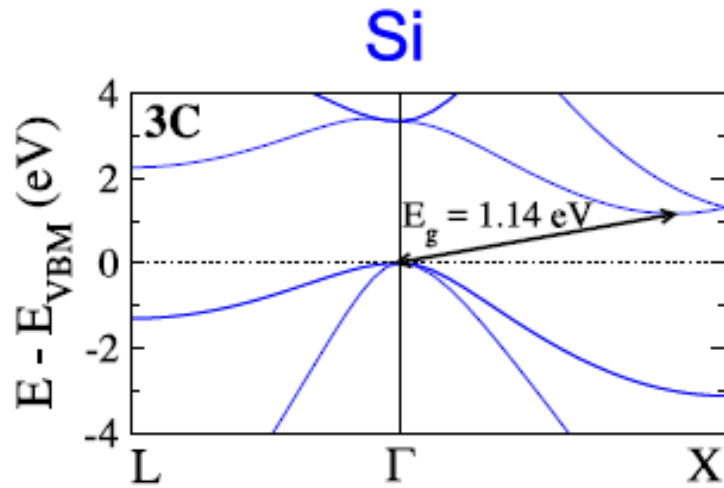
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# Band structure of 2H and 3C bulk

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DFT – HSE06



# Band gaps of 2H and 3C bulk



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Table 2: The calculated band gaps of Si and Ge.

Element	Crystal structure	$E_g$ at $\Gamma$ (eV)	$E_g$ (eV)	$E_g$ (SOC) (eV)	Band-edge	Notes
Si	Cubic-diamond (3C)	3.29	1.14		$\Gamma - X$	This work (HSE06)
		3.20	1.14		$\Gamma - X$	Cal. Ref. <sup>16</sup>
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T. Kaewmaraya et al., J. Phys. Chem. C 121, 5820 (2017)

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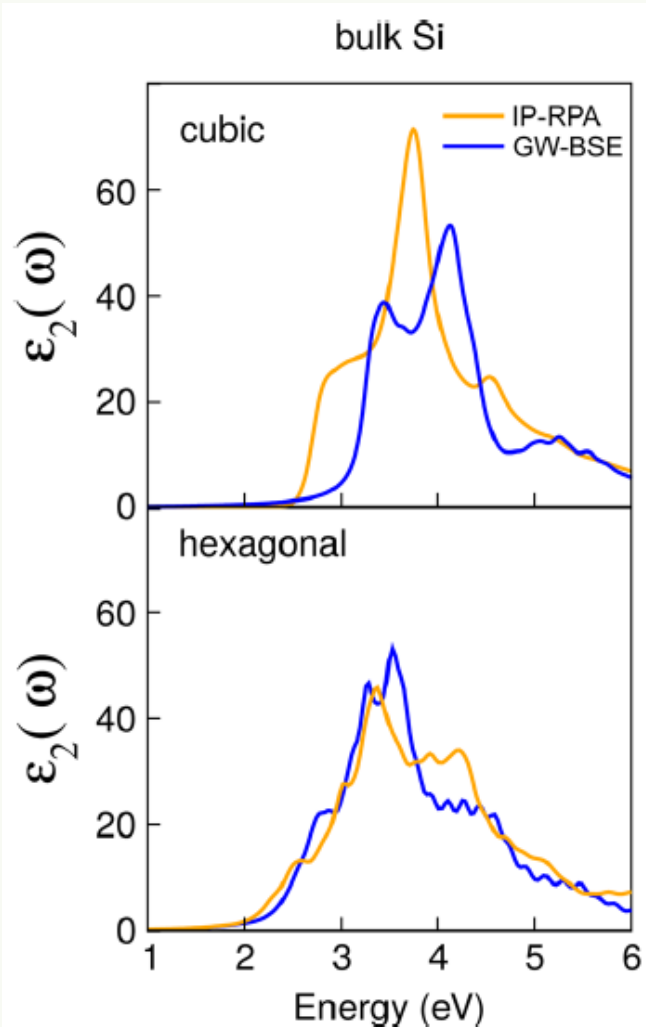
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# Optical spectra of 2H and 3C Si bulk

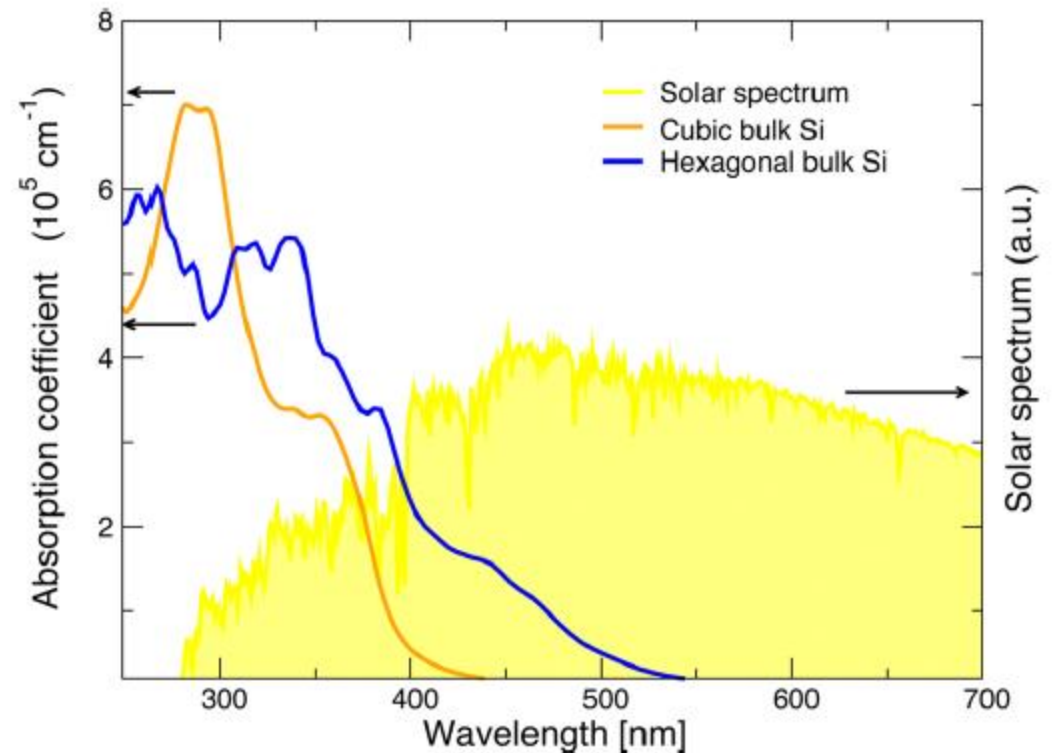
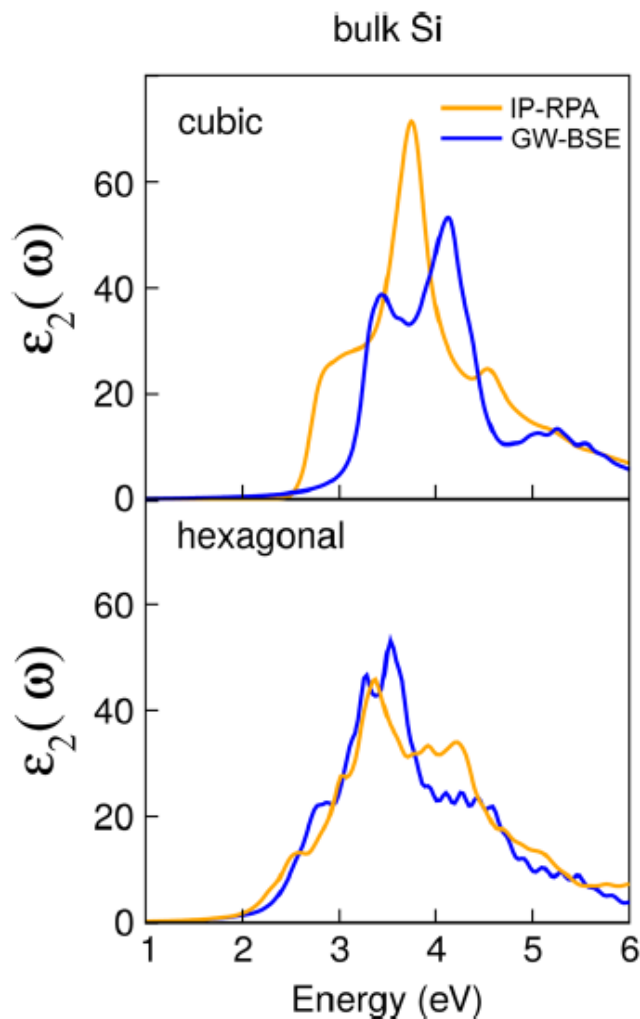
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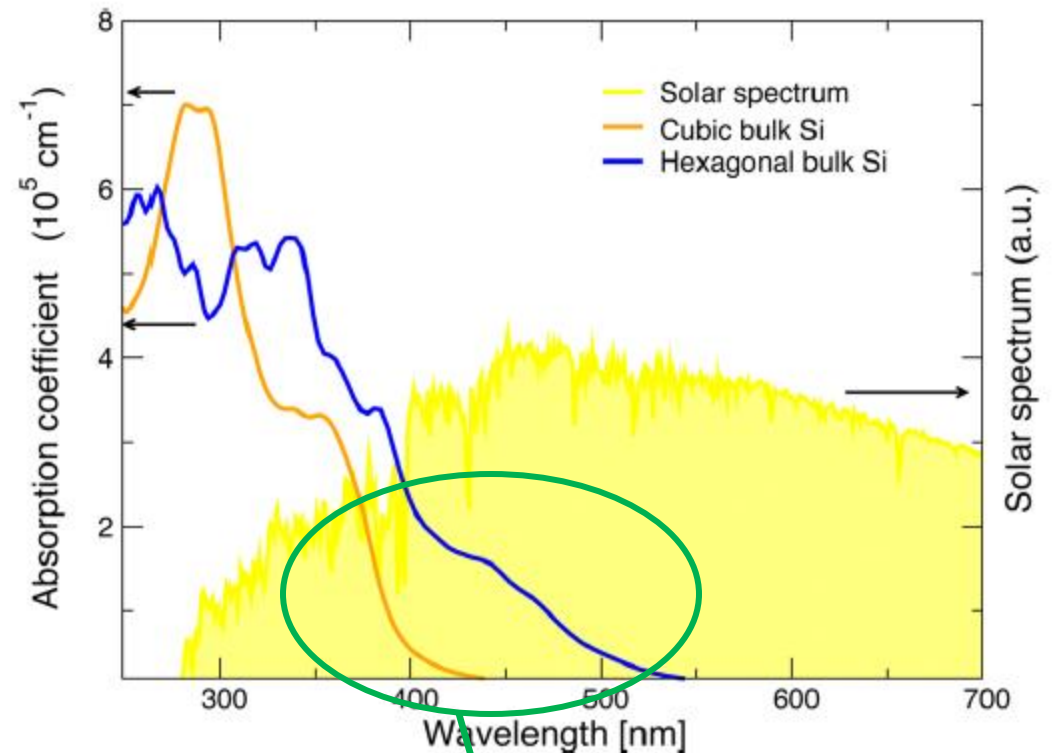
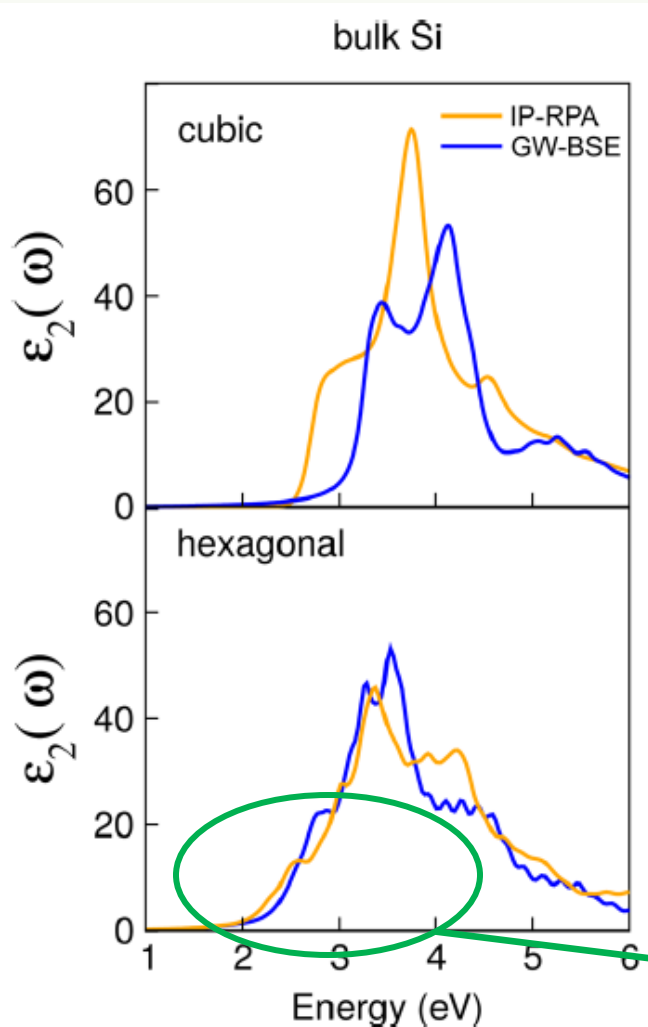
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GW – BSE



2H bulk phase

Larger absorption in the visible

Bigger overlap close to the max of solar spectrum

# Some comments



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**What happens when we reduce the size?**  
**Should we expect the same electronic and optical behaviour?**

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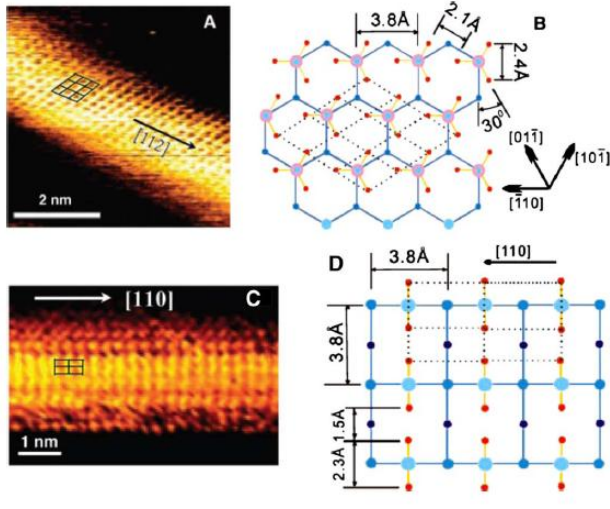
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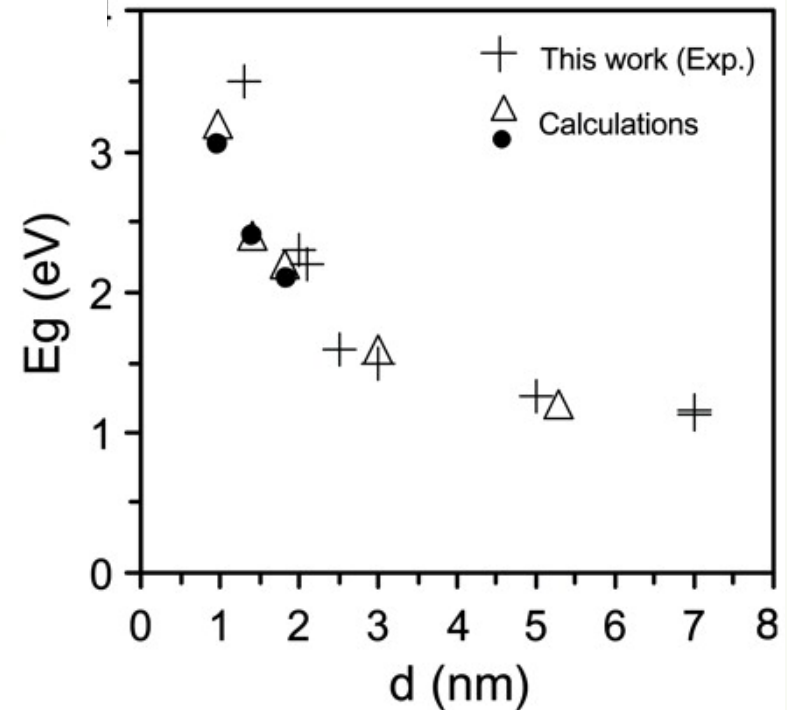
# Quantum confinement in 3C Si NWs

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Ma, D. D. D. et al., Science 299, 1874 (2003)

## STS measurements



Ma, D. D. D. et al., Science 299, 1874 (2003)

$$E_{\text{gap}} = E_{\text{bulk, gap}} + C(1/d)^{\alpha}$$

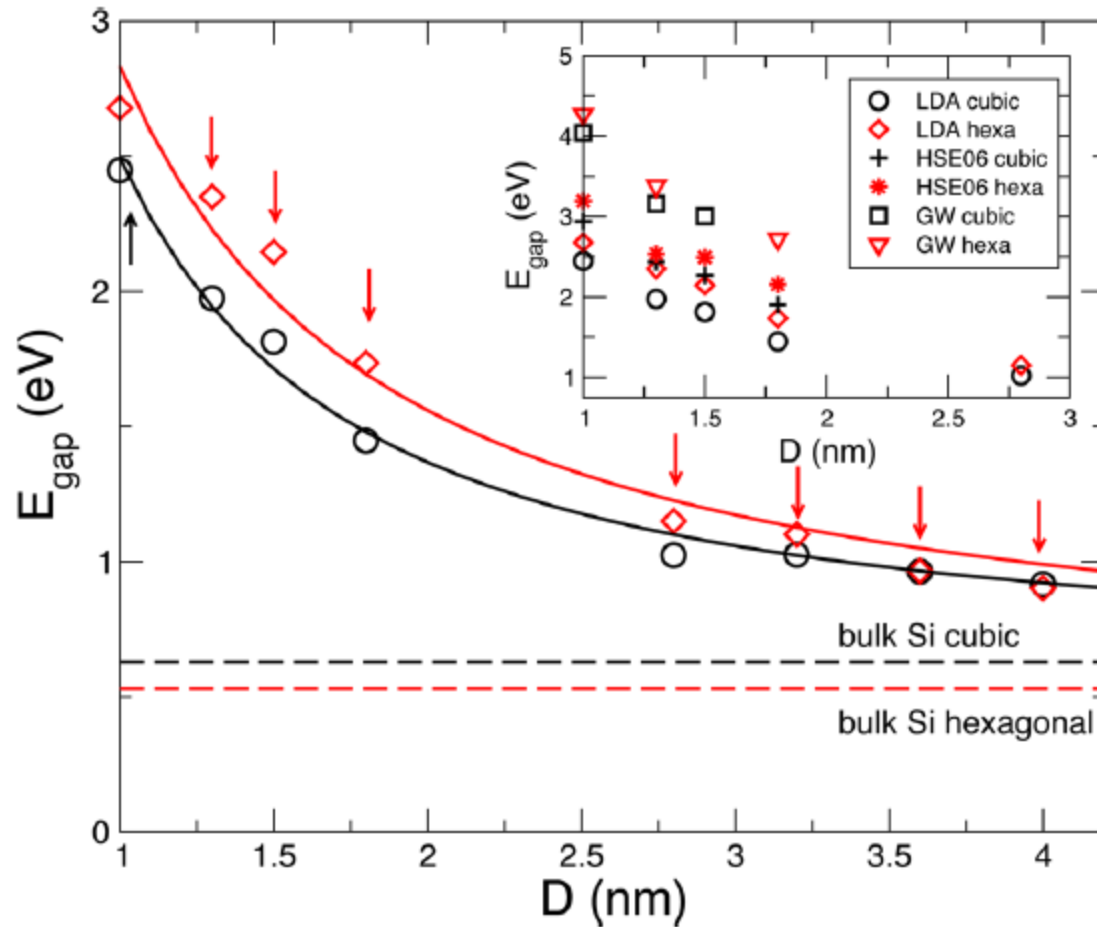
with  $\alpha=0.9 - 1.1$

M. Bruno et al., Phys. Rev. Lett. 98, 036807 (2007)

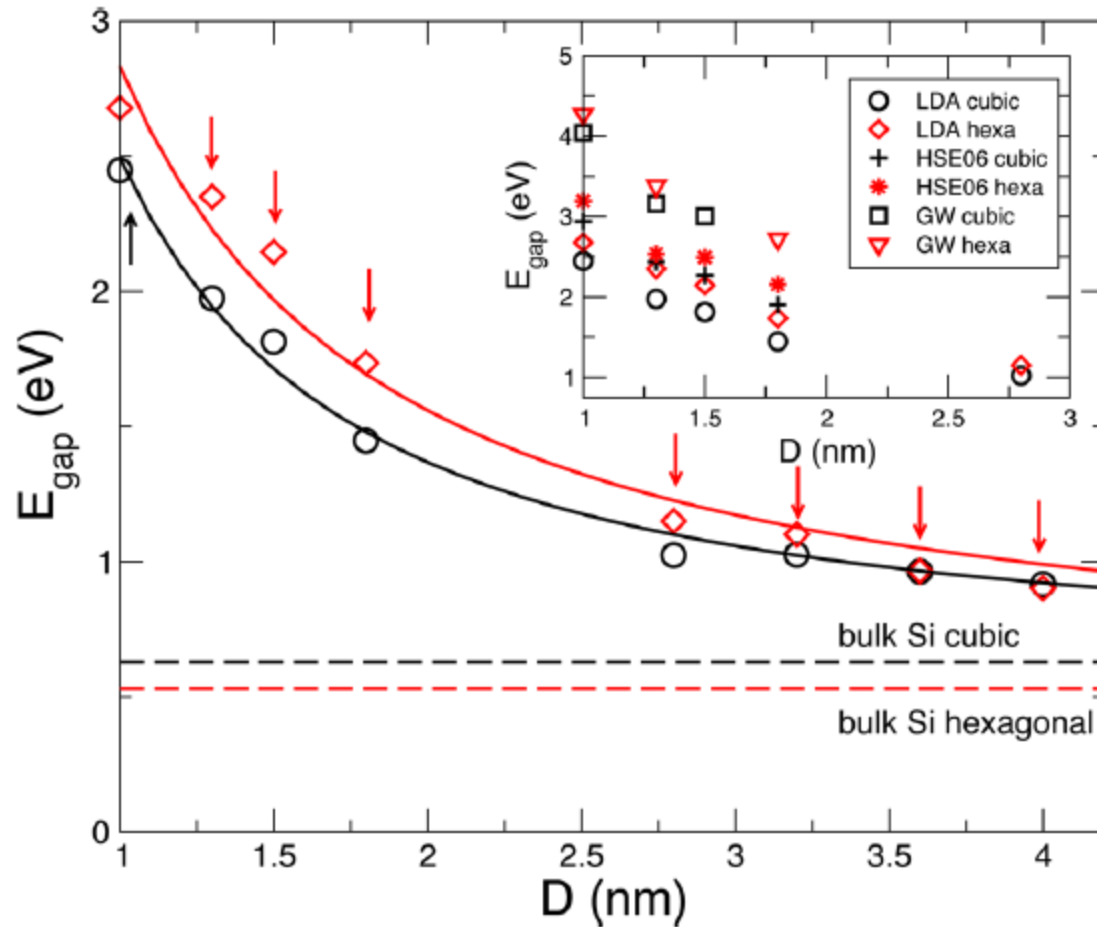


# Band gap of 3C and 2H Si NWs

# Band gap of 3C and 2H Si NWs



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DFT – LDA

$$\alpha_{\text{hex}} = 1.16$$

$$\alpha_{\text{cub}} = 1.34$$

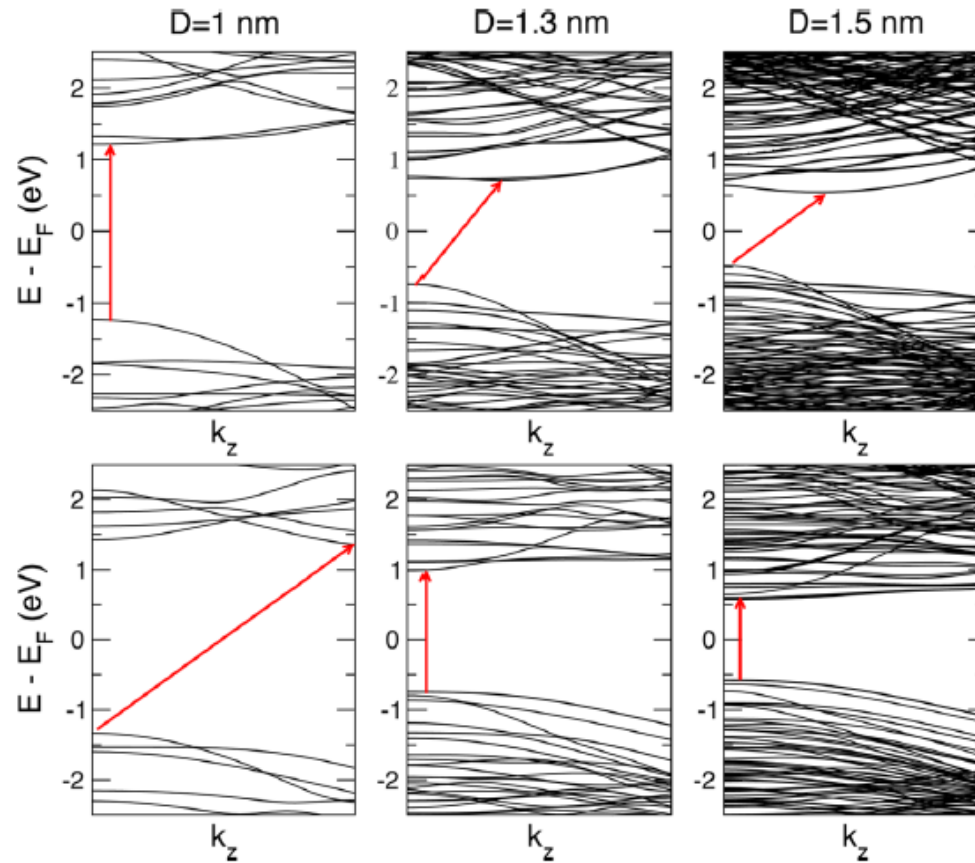
$$E_{\text{gap}} = E_{\text{bulk, gap}} + C(1/d)^{\alpha}$$

# Bands and optical absorption

## DFT – LDA

3C

2H



M. Amato et al., Nano Lett. 16, 5694 (2016)

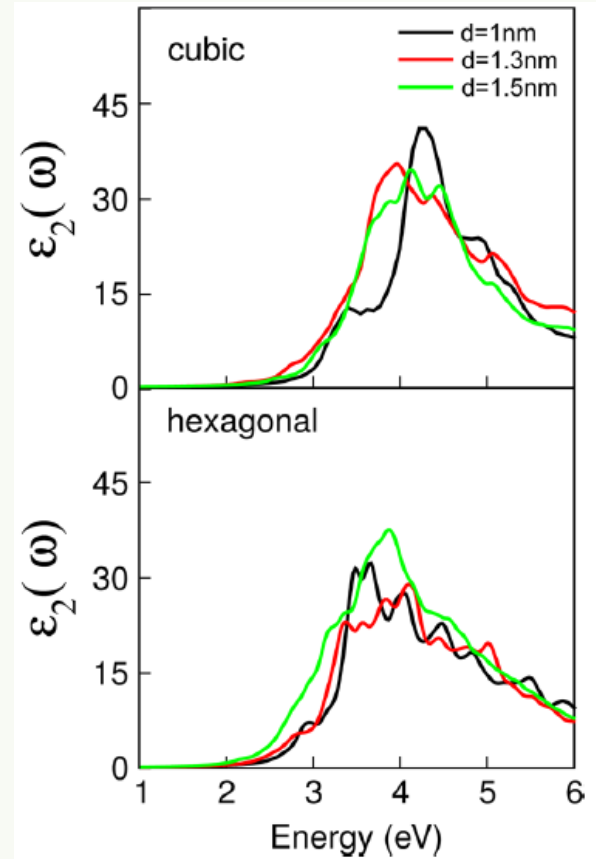
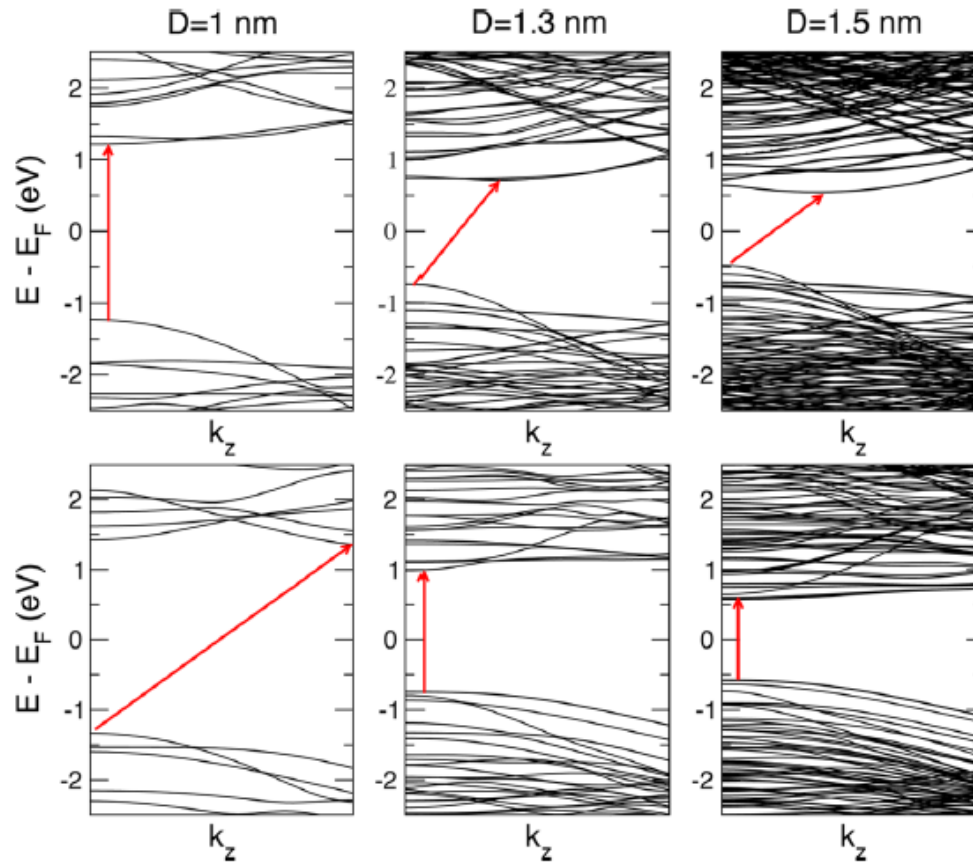
# Bands and optical absorption

DFT – LDA

IP – RPA

3C

2H



M. Amato et al., Nano Lett. 16, 5694 (2016)

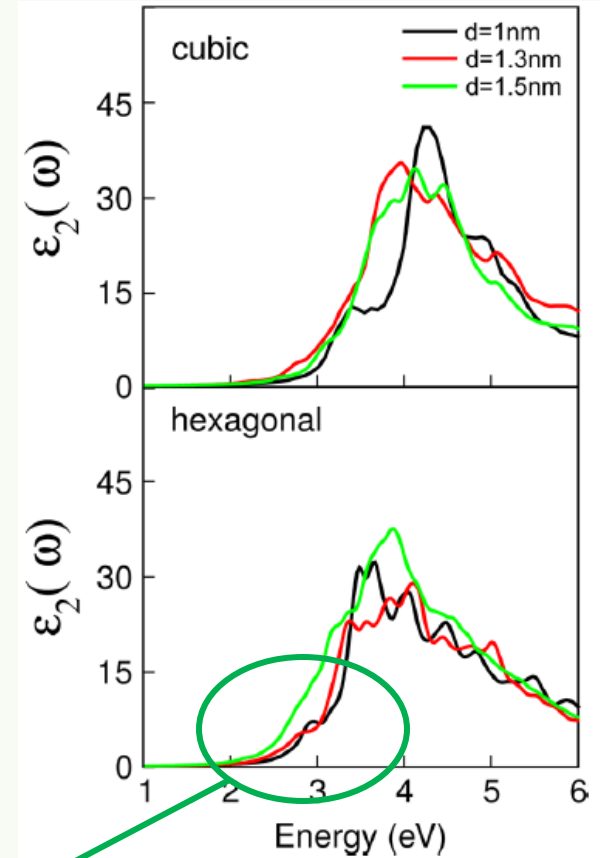
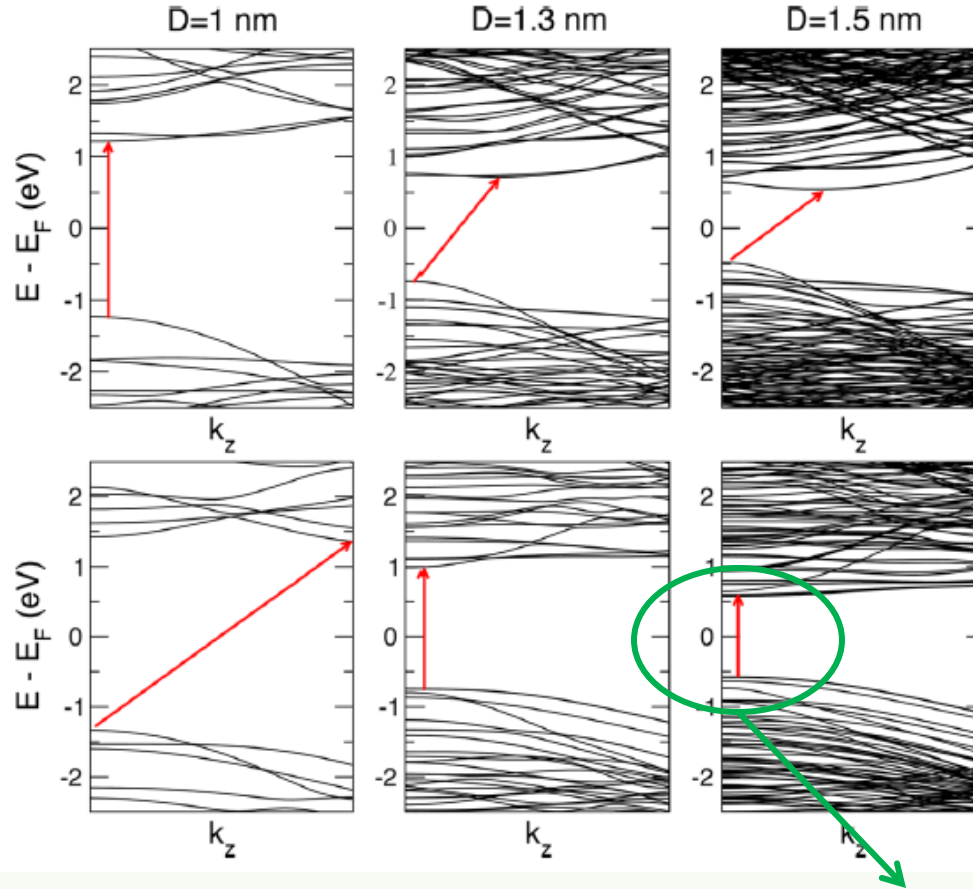
# Bands and optical absorption

DFT – LDA

IP – RPA

3C

2H



M. Amato et al., Nano Lett. 16, 5694 (2016)

2H <111> Si NWs

Larger absorption in the visible

# Outline

- Introduction on nanoscience
- Ab initio simulations: Density Functional Theory (DFT)
- Quantum confinement in Si nanowires (Si NWs)
  - Electronic structure
  - Transport properties
  - Doping effects
- Crystal phase engineering in Si NWs**
  - Crystal phase effects on bulk
  - Crystal phase effects on NWs**
  - Homojunctions
- Conclusions

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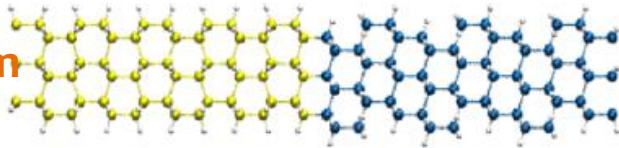


# Band offset in 3C/2H NWs junctions

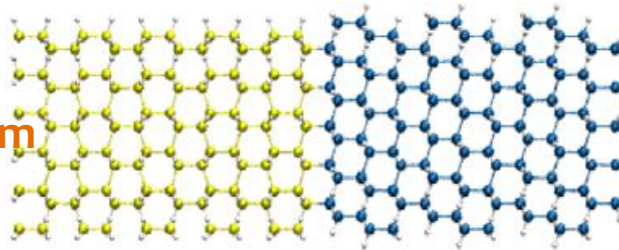
# Band offset in $3C/2H$ NWs junctions

$[0001]_{\text{hex}}/\langle 111 \rangle_{\text{cub}}$  Si NWs

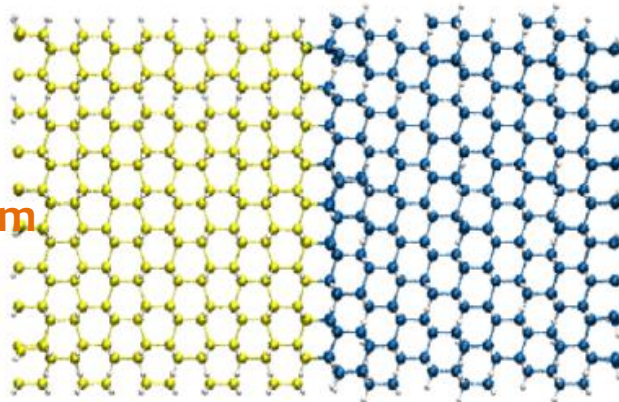
D=1 nm



D=1.8 nm



D=3.2 nm



ABABABABABCBBCABBCAB

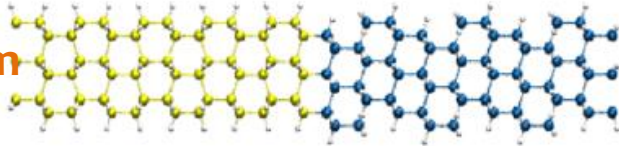
$L_{2H} = \text{constant}$

$L_{3C} = \text{constant}$

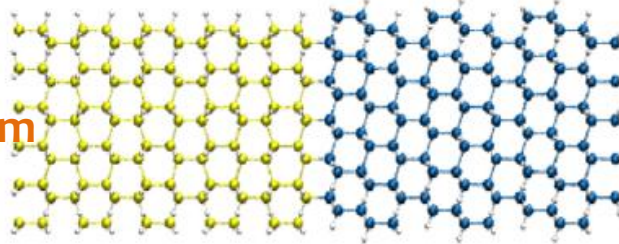
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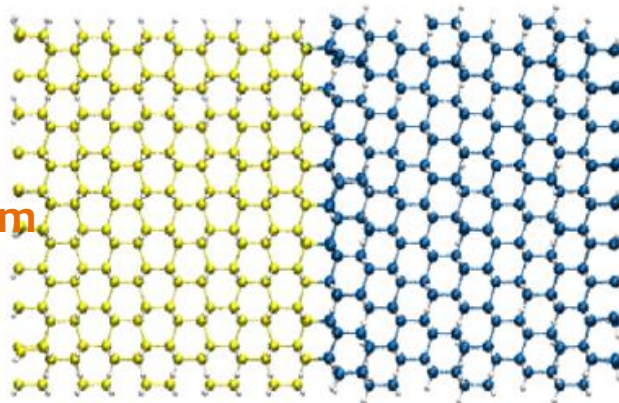
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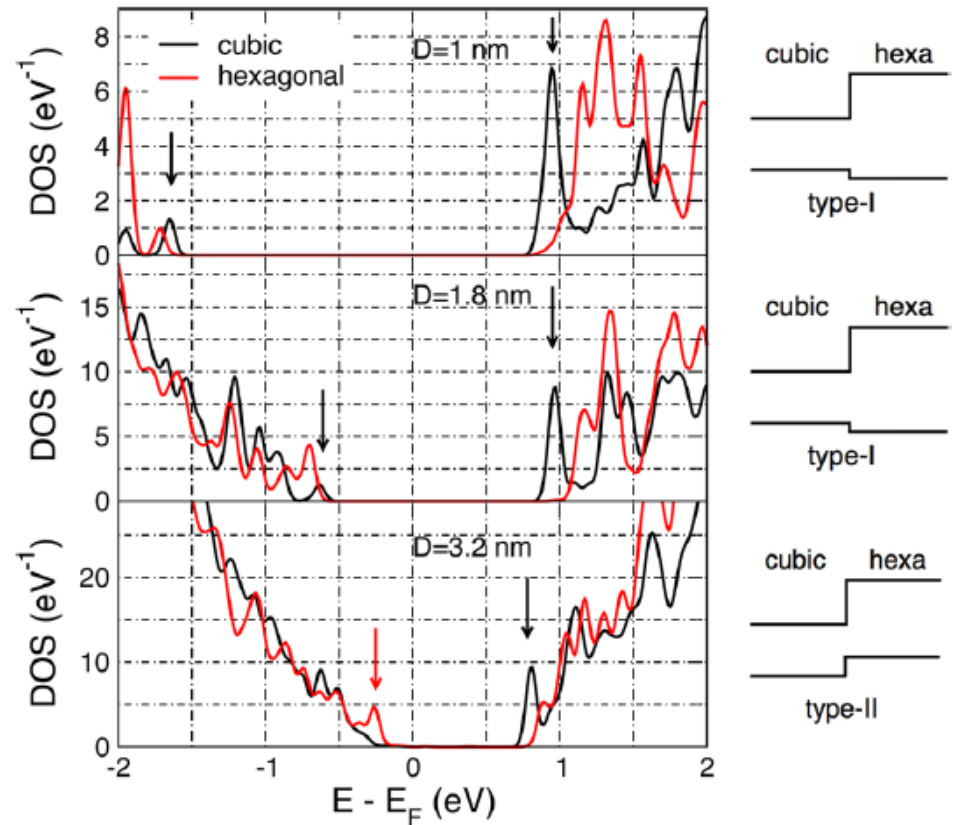


ABABABABABCBABCABCB

$L_{2\text{H}} = \text{constant}$

$L_{3\text{C}} = \text{constant}$

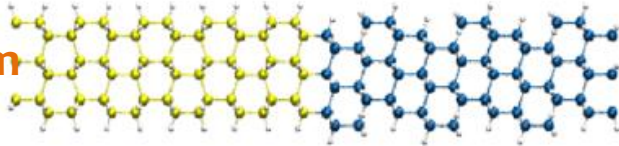
Projected Density of States (PDOS)



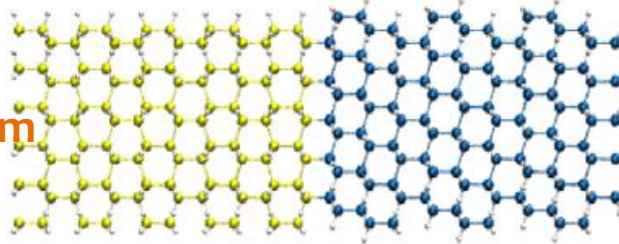
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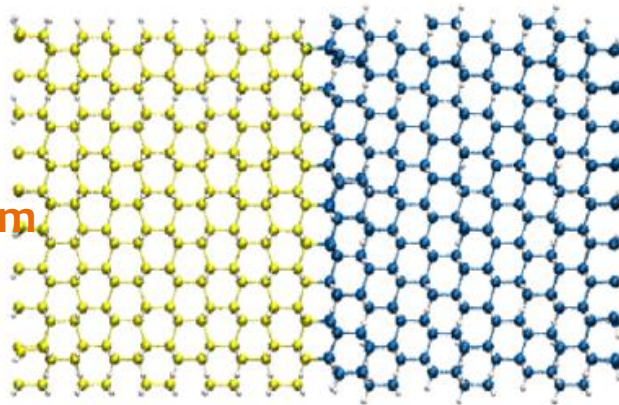
D=1 nm



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D=3.2 nm

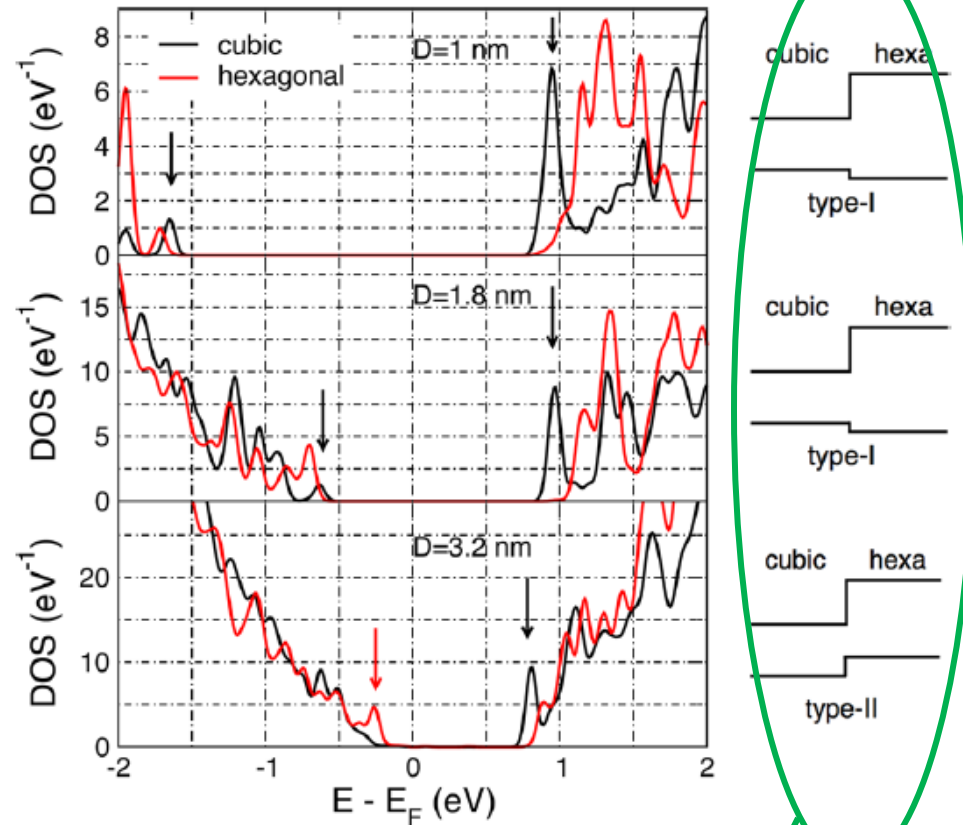


ABABABABABABCABBCAB

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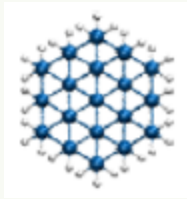
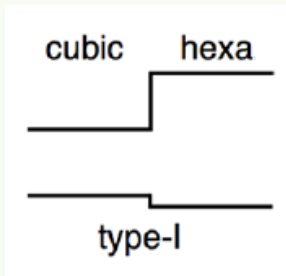
Switch of the band offset

Different confinement of 2H and 3C NWs

# Modulation of the band offset with D

# Modulation of the band offset with D

$D=1\text{ nm}$

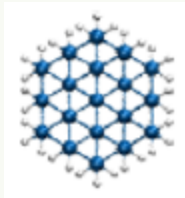
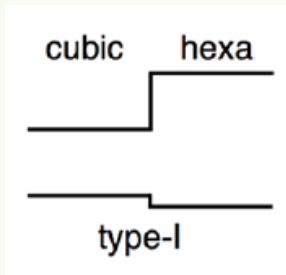


*Keeping the junction length constant*

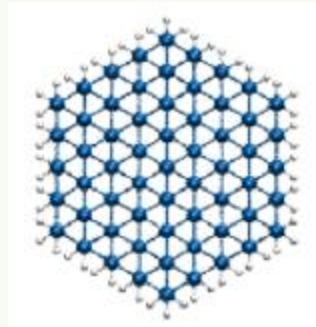
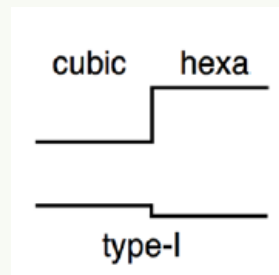
**D**

# Modulation of the band offset with D

D=1 nm



D=1.8 nm



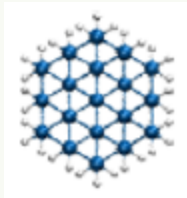
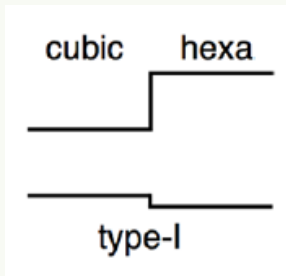
Keeping the junction length constant

D

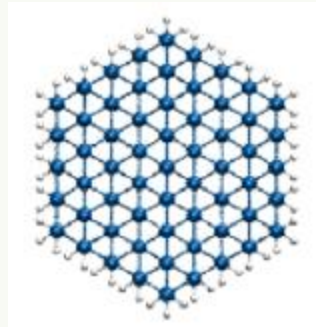
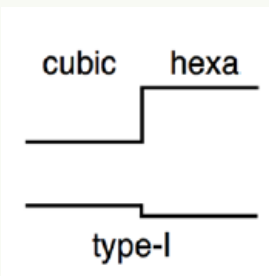


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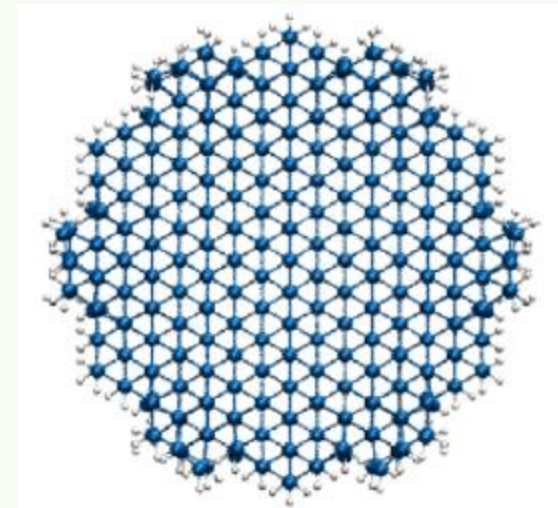
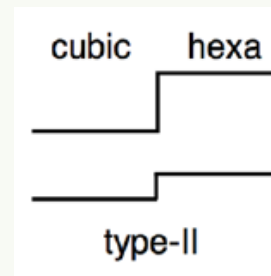
D=1 nm



D=1.8 nm



D=3.2 nm



Keeping the junction length constant

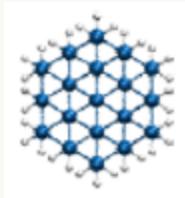
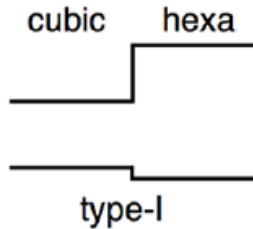
D



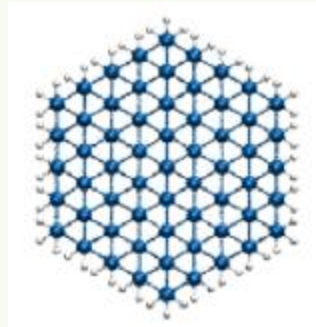
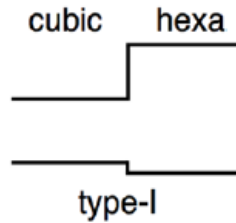
# Modulation of the band offset with D

2H segments act  
as a quantum dot

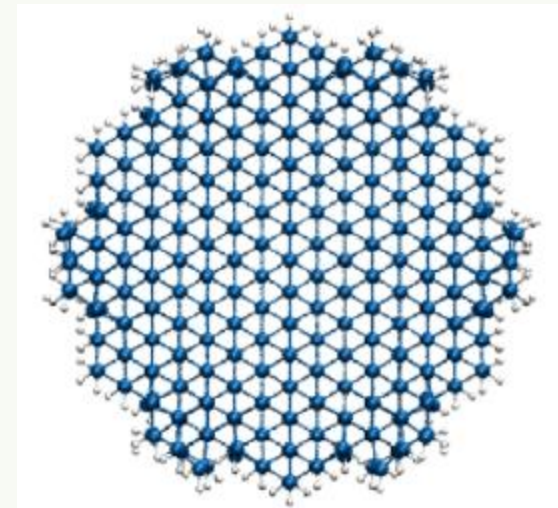
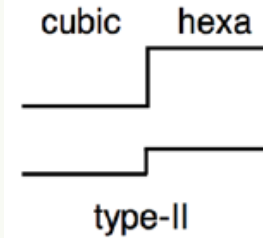
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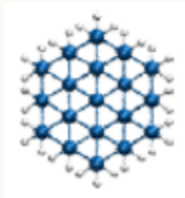
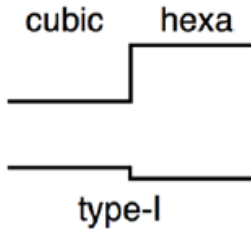
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D

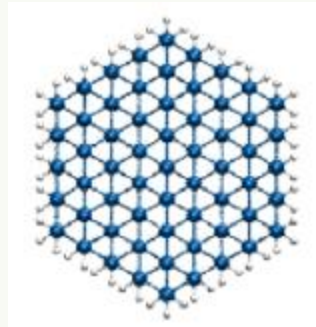
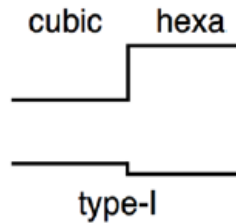
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D=1 nm

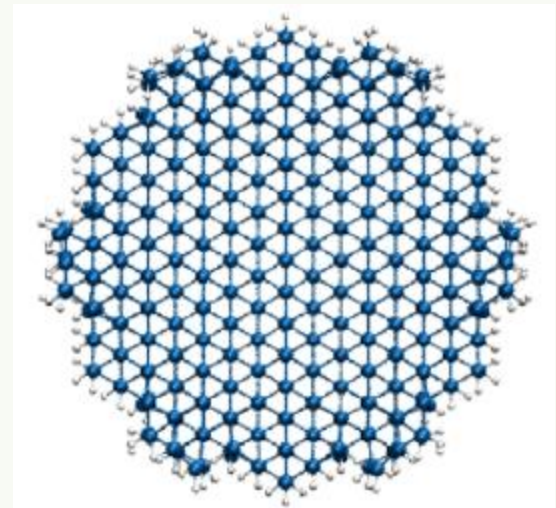
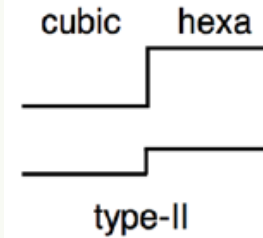


D=1.8 nm



2H segments induce  
e-h separation

D=3.2 nm



Keeping the junction length constant

D

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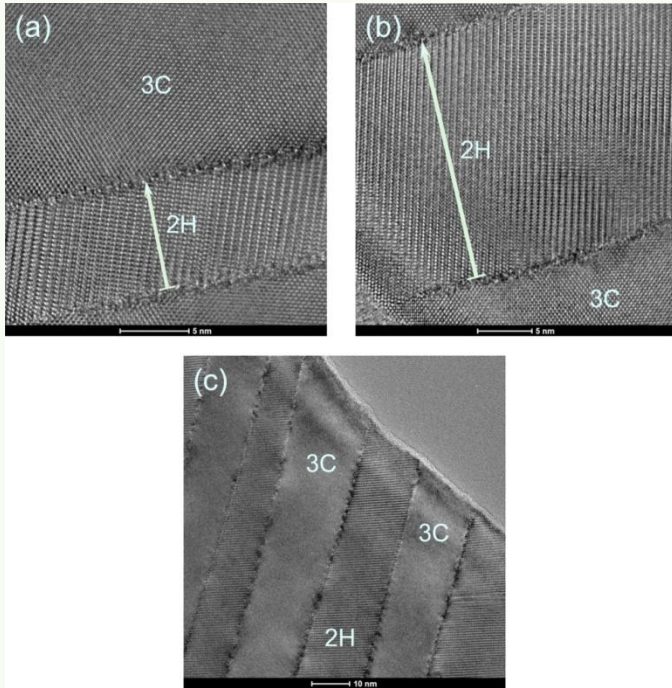
**Can we say something on thicker wires?**

# $3C/2H$ NWs junctions with large D



# 3C/2H NWs junctions with large D

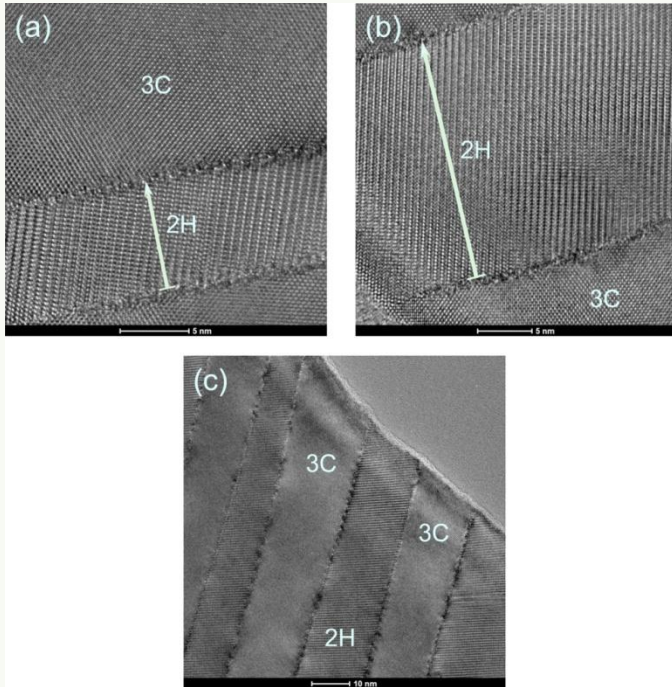
## Modulation of the 2H region width with T



# $3C/2H$ NWs junctions with large D

Theoretical  
Model

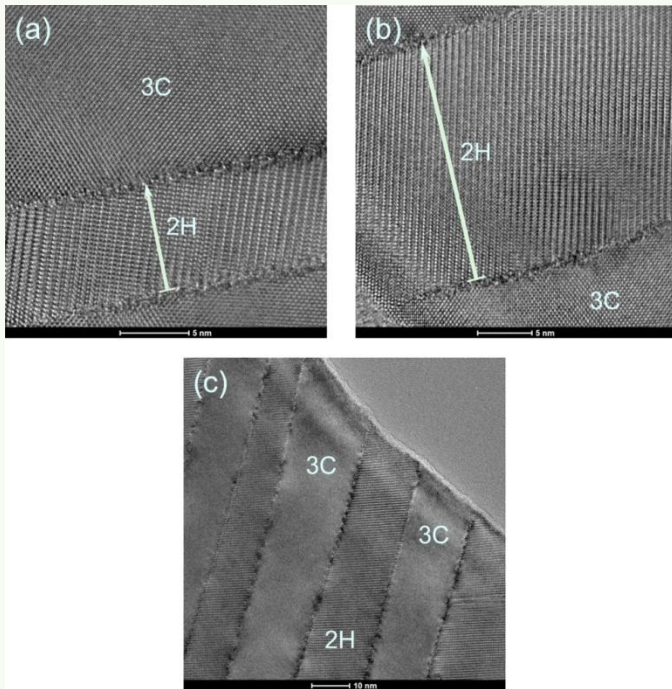
Modulation of the 2H region width with T



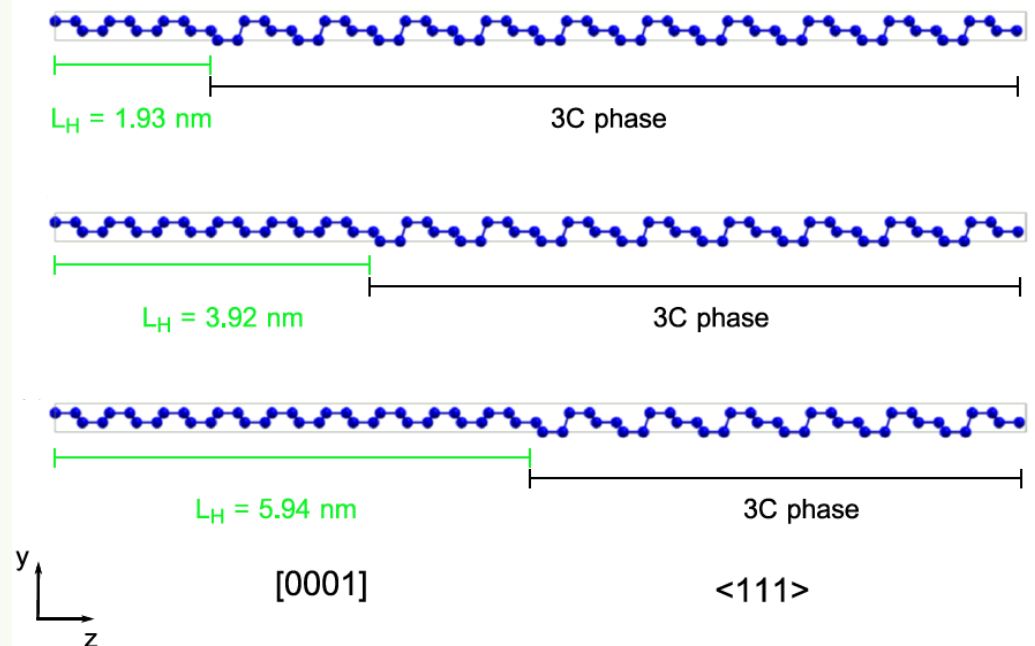
# 3C/2H NWs junctions with large D

Theoretical  
Model

Modulation of the 2H region width with T



Unit cell of  $\langle 111 \rangle_{\text{cub}} / [0001]_{\text{hex}}$  Si and Ge homojunctions  
Large diameters systems ~ bulk junctions

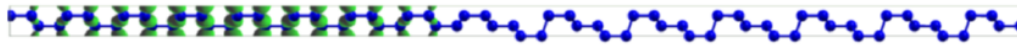


# Band offset in larger wires

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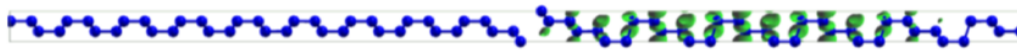
## Si NW junctions

VBM



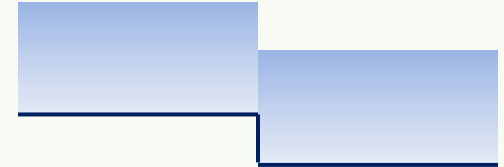
2H phase

CBM



3C phase

## Type II

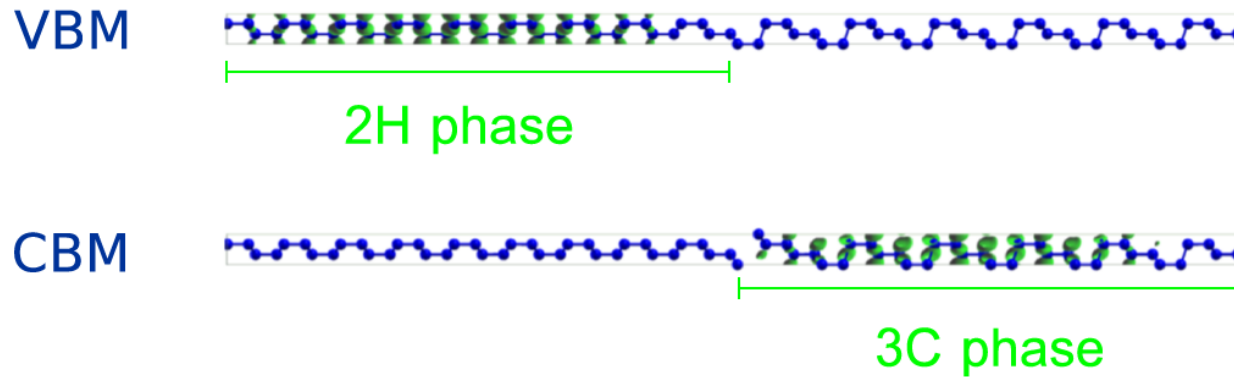


2H phase

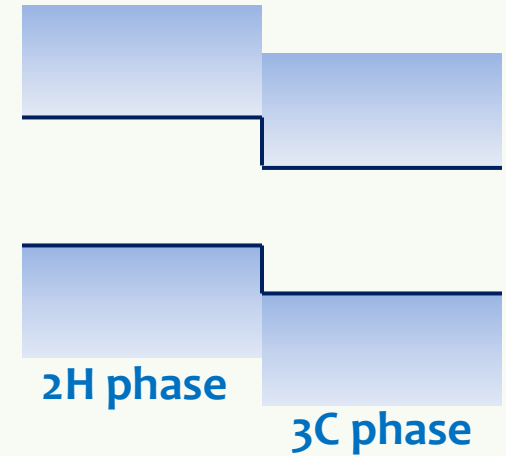
3C phase

# Band offset in larger wires

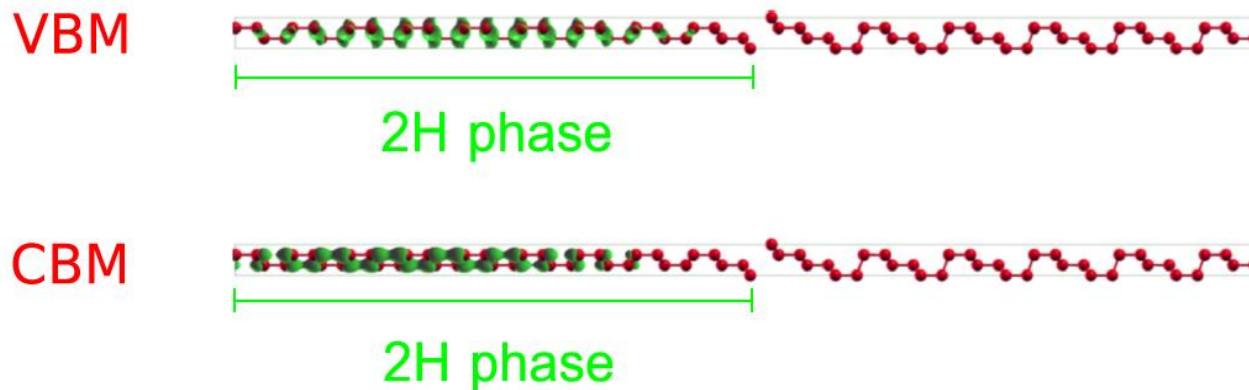
## Si NW junctions



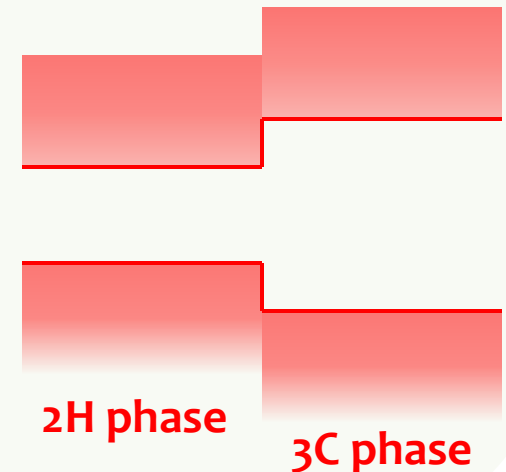
## Type II



## Ge NW junctions



## Type I



# Some comments

# Some comments

- ▶ **Si junctions** are characterized by a **type-II band alignment**, with valence and conduction band edges spatially separated could be of some benefit for **photovoltaics** as well as other advanced technological applications
- ▶ Moreover, in Si homojunctions varying the **thickness of 2H** region **keep almost unaltered** the **band alignment** nature, the band gap values and the offset magnitudes



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